

# Gamow-Teller strength distributions at finite temperatures and electron capture rates in stellar environment

A. Dzhihev<sup>1,2</sup>, A. Vdovin<sup>1</sup>, V. Ponomarev<sup>3</sup>, J. Wambach<sup>3,4</sup>,  
K. Langanke<sup>3,4,5</sup>, and G. Martínez-Pinedo<sup>4</sup>

<sup>1</sup>Bogoliubov Laboratory of Theoretical Physics, JINR, Dubna, Russia

<sup>2</sup>Université Libre de Bruxelles, Belgium

<sup>3</sup>Institut für Kernphysik, TU Darmstadt, Germany

<sup>4</sup>GSI, Darmstadt, Germany

<sup>5</sup>Frankfurt Institute for Advanced Studies, Germany

**11th Symposium on Nuclei in the Cosmos (NIC XI)**

Heidelberg, Germany, July 19 – 23, 2010

## Electron capture on nuclei during the core collapse



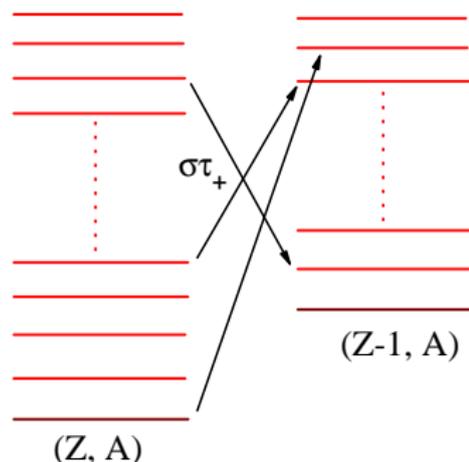
- Reduces the number of electrons ( $Y_e$ ) and keeps the core cool;
- Decreases Chandrasekar mass ( $M_{Ch} \approx 1.44(2Y_e)^2 M_\odot$ );
- Dominated by Gamow-Teller transitions caused by the  $\sigma\tau_+$  operator.

In stellar environments EC takes place at finite temperatures  $T = 0.2 - 2.0$  MeV ( $0.1 \text{ MeV} \approx 1.2 \times 10^9 \text{ K}$ ) and  $GT_+$  distributions for nuclear **excited states** are needed.

EC rates are computed by:

- Shell-Model for  $A < 65$ ;
- Hybrid model (SMMC+RPA) for  $A > 65$ .

$GT_+$  transitions at finite temperatures



## Basic ingredients of the Thermal QRPA (pnQRPA + TFD) approach:

- Thermo Field Dynamics
  - Thermal Hamiltonian:  $\mathcal{H} = H - \tilde{H}$ , where  $H|n\rangle = E_n|n\rangle$  and  $\tilde{H}|\tilde{n}\rangle = E_n|\tilde{n}\rangle$ ;
  - Thermal vacuum:  $\mathcal{H}|0(T)\rangle = 0$  and  $\ll A \gg = \langle 0(T)|A|0(T)\rangle$ ;
- The QPM Hamiltonian:  $H = H_{W.S.} + H_{BCS} + H_{ph}$ .

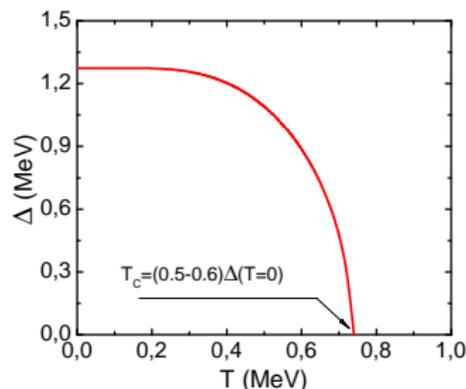
## Diagonalization of $\mathcal{H}$ within TQRPA:

- Thermal quasiparticles:

$$\mathcal{H}_{W.S.+BCS} \approx \sum_{i,\tau=p,n} \epsilon_{i\tau} (\beta_{i\tau}^\dagger \beta_{i\tau} - \tilde{\beta}_{i\tau}^\dagger \tilde{\beta}_{i\tau});$$

- Thermal phonons:

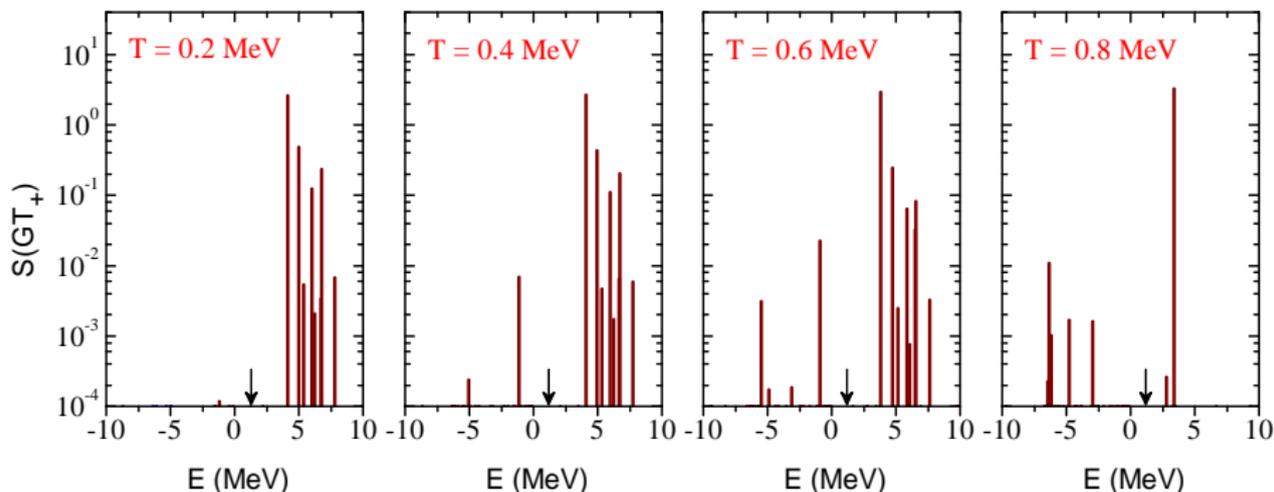
$$\mathcal{H} \approx \sum_{k\lambda} \omega_{k\lambda} (\mathcal{Q}_{k\lambda}^\dagger \mathcal{Q}_{k\lambda} - \tilde{\mathcal{Q}}_{k\lambda}^\dagger \tilde{\mathcal{Q}}_{k\lambda})$$



Transition strength:  $S_{k\lambda} = \langle \mathcal{Q}_{k\lambda}^\dagger | \hat{\mathcal{O}}_\lambda | 0(T) \rangle$  and  $S_{\tilde{k}\lambda} = \langle \tilde{\mathcal{Q}}_{k\lambda}^\dagger | \hat{\mathcal{O}}_\lambda | 0(T) \rangle$ .

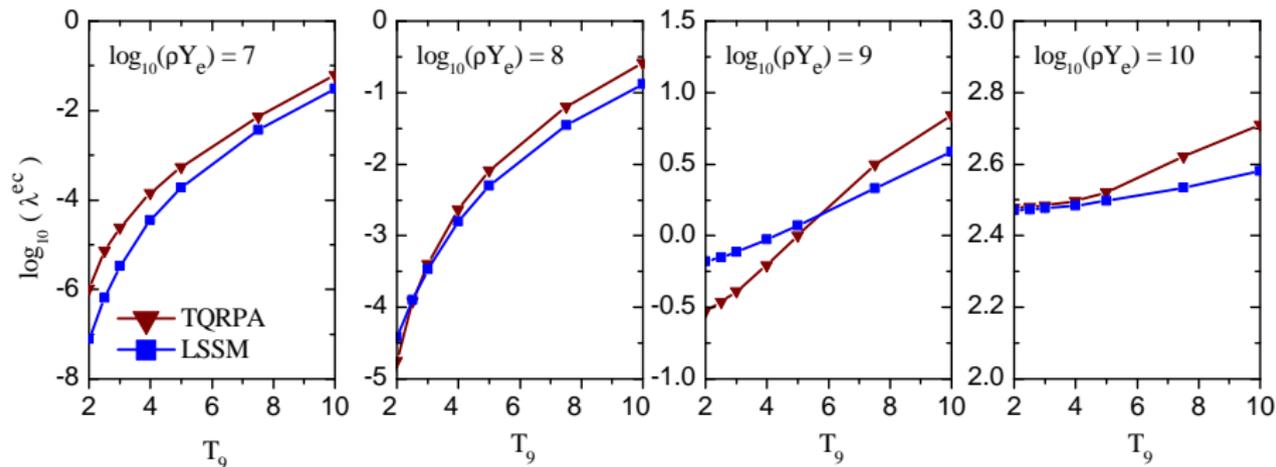
## Strength distributions of $GT_+$ transitions in $^{54}\text{Fe}$

Iron group nuclei are essential at the early presupernova collapse,  
 $T = 0.2 - 0.8 \text{ MeV}$  and  $\rho = 10^7 - 10^{10} \text{ g cm}^{-3}$ .



The black arrows indicate the zero-temperature EC threshold  
 $Q = M(^{54}\text{Fe}) - M(^{54}\text{Mg}) = 1.21 \text{ MeV}$ .

# Electron capture rates for $^{54}\text{Fe}$

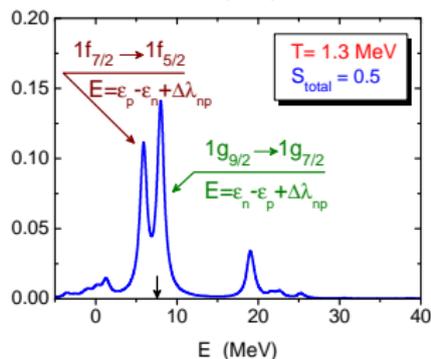
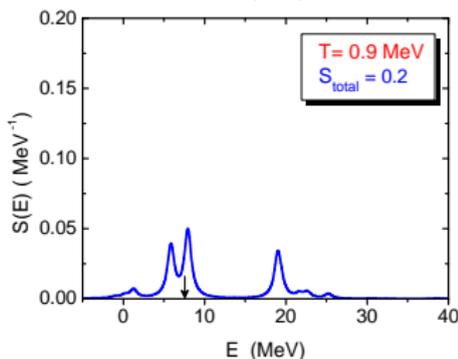
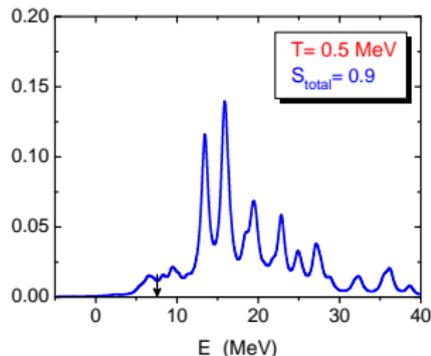
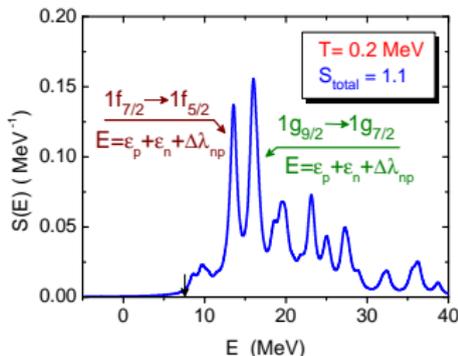


$$\lambda^{ec} = \frac{\ln 2}{6150\text{s}} \sum_i S_i(GT_+) F_i,$$

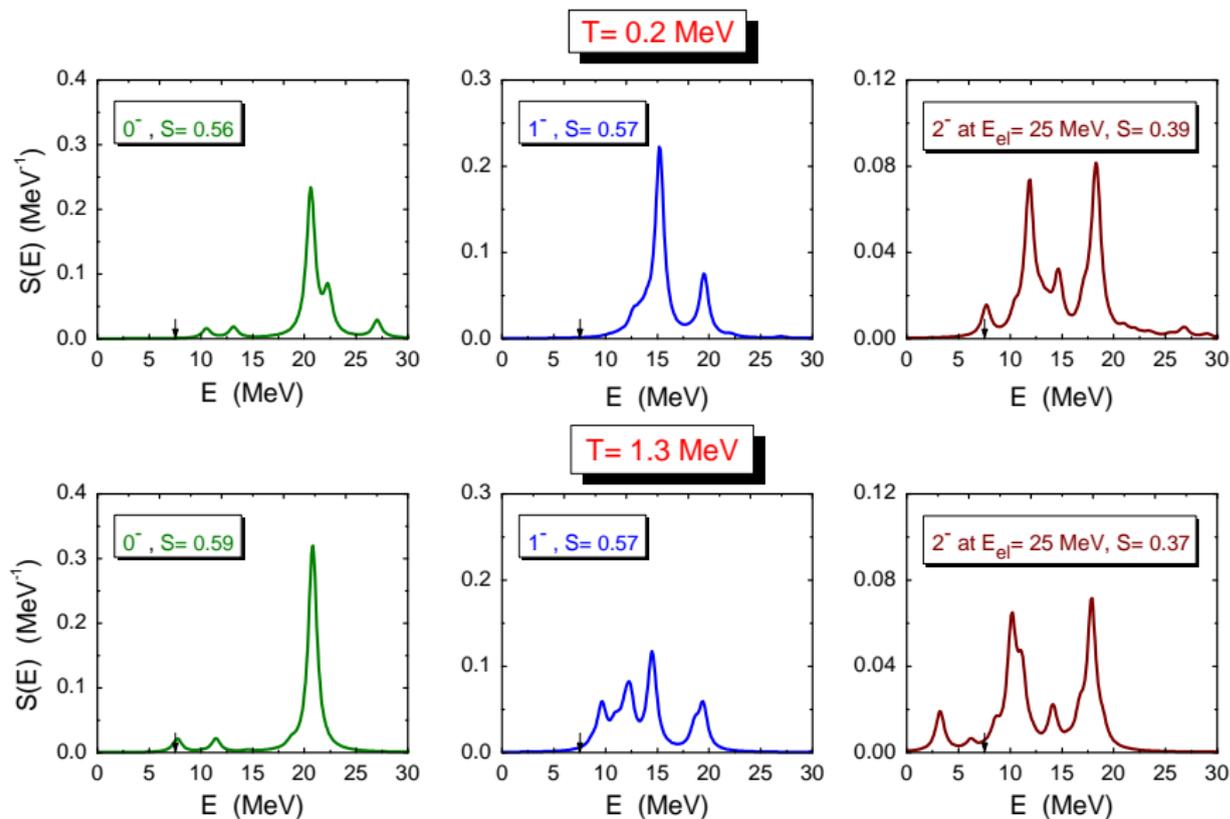
$T_9 = 10^9$  K; density  $\rho Y_e$  in  $\text{g cm}^{-3}$ .

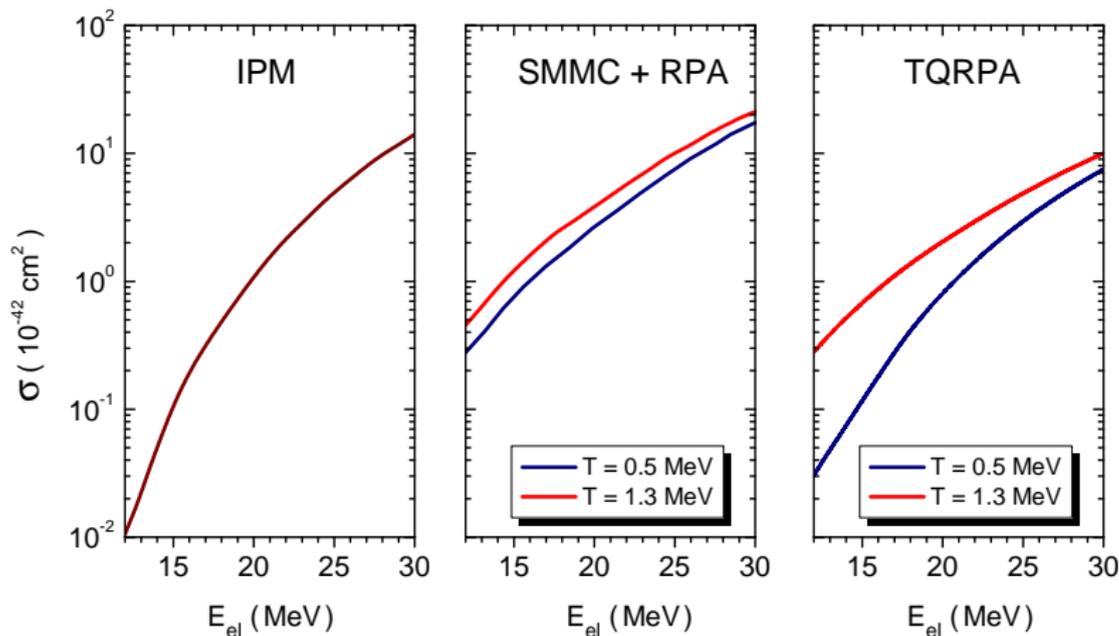
# Strength distributions of $GT_+$ transitions in $^{76}\text{Ge}$

- Neutron-rich nuclei with  $N \geq 40$  and  $Z \leq 40$  dominate the nuclear composition for  $\rho > 10^{10} \text{ g cm}^{-3}$  and  $T > 0.8 \text{ MeV}$
- Unblocking mechanisms: thermal excitations and configuration mixing.



# Strength distributions of first-forbidden $p \rightarrow n$ transitions in $^{76}\text{Ge}$





$$\sigma(E_{el}, T) = \frac{G_w}{2\pi} F(Z, E_{el}) \sum_{\lambda i} (E_{el} - E_{\lambda i})^2 S_{\lambda i}, \quad \lambda = 1^-, 0^+, 1^+, 2^+.$$

- The novel approach to study thermal effects on the  $GT_+$  strength distributions and electron captures on nuclei in stellar environment have been presented. It was shown that thermal effects shift  $GT_+$  centroid to lower excitation energies and make possible negative- and low-energy transitions. It was found that the unblocking effect for  $GT_+$  transitions in neutron-rich nuclei is more sensitive to increasing temperature than it was predicted by the hybrid model.
- To improve the predictive power of the approach it is desirable to combine our TFD-based method with self-consistent QRPA calculations based on more realistic effective interactions. Another direction is the inclusion of correlations beyond TQRPA by coupling one-phonon states with more complex configurations like it was done at zero temperature within the QPM.



**Acknowledgments:** *This work is supported by the Heisenberg-Landau Program, the DFG grant (SFB 634), the Helmholtz Alliance EMMI and HIC for Fair.*