

The Nuclear Equation of State: What Do Neutron Stars Tell Us?

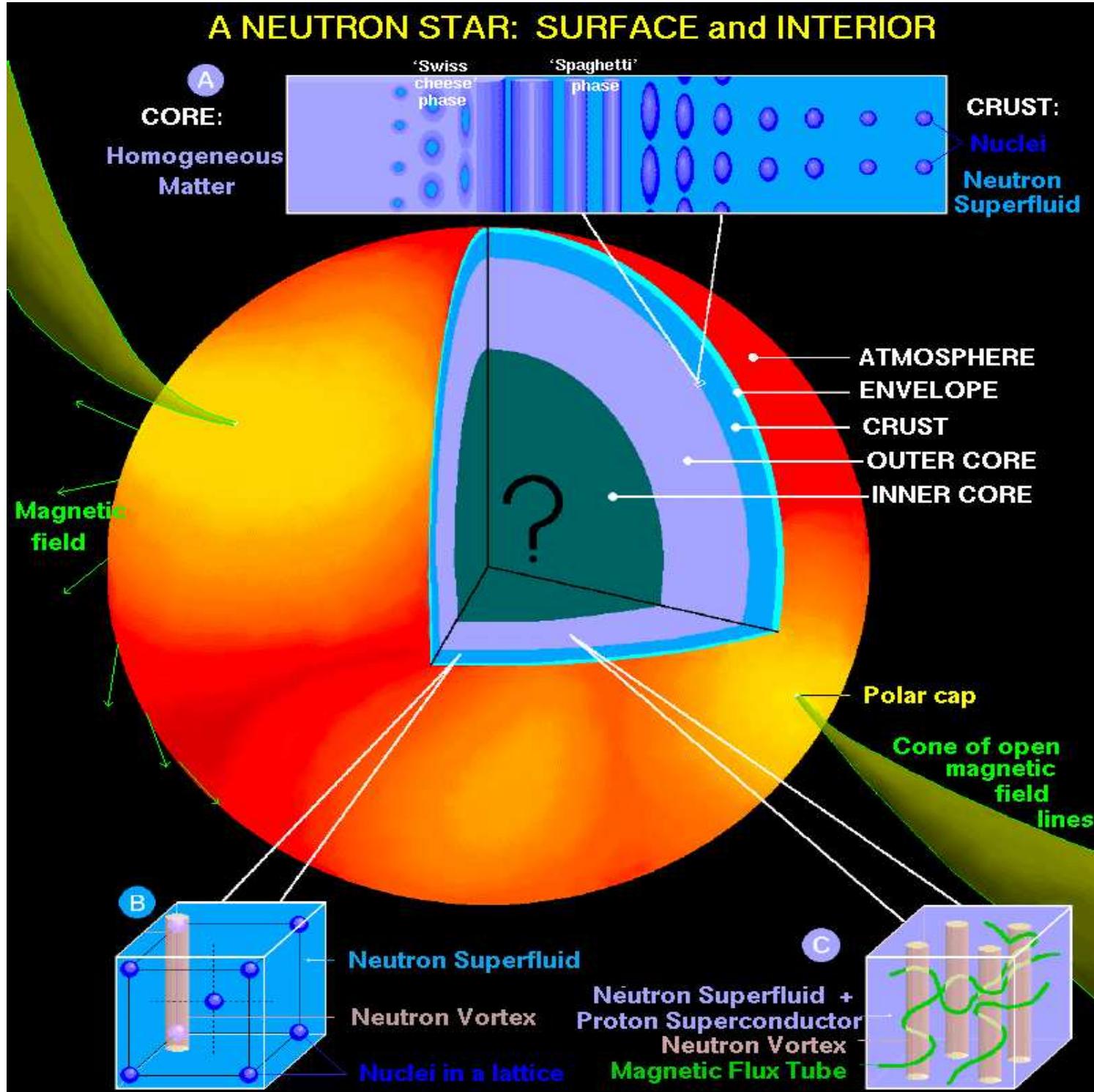
James M. Lattimer

Department of Physics & Astronomy
Stony Brook University

Collaborators:

- Ed Brown (MSU)
- Kai Hebeler (TRIUMF)
- Bryan Kim (Stony Brook)
- Ralph Neuhäuser (Jena)
- Chris Pethick (NORDITA)
- Madappa Prakash (Ohio)
- Achim Schwenk (GSI)
- Andrew Steiner (MSU)
- Fred Walter (Stony Brook)

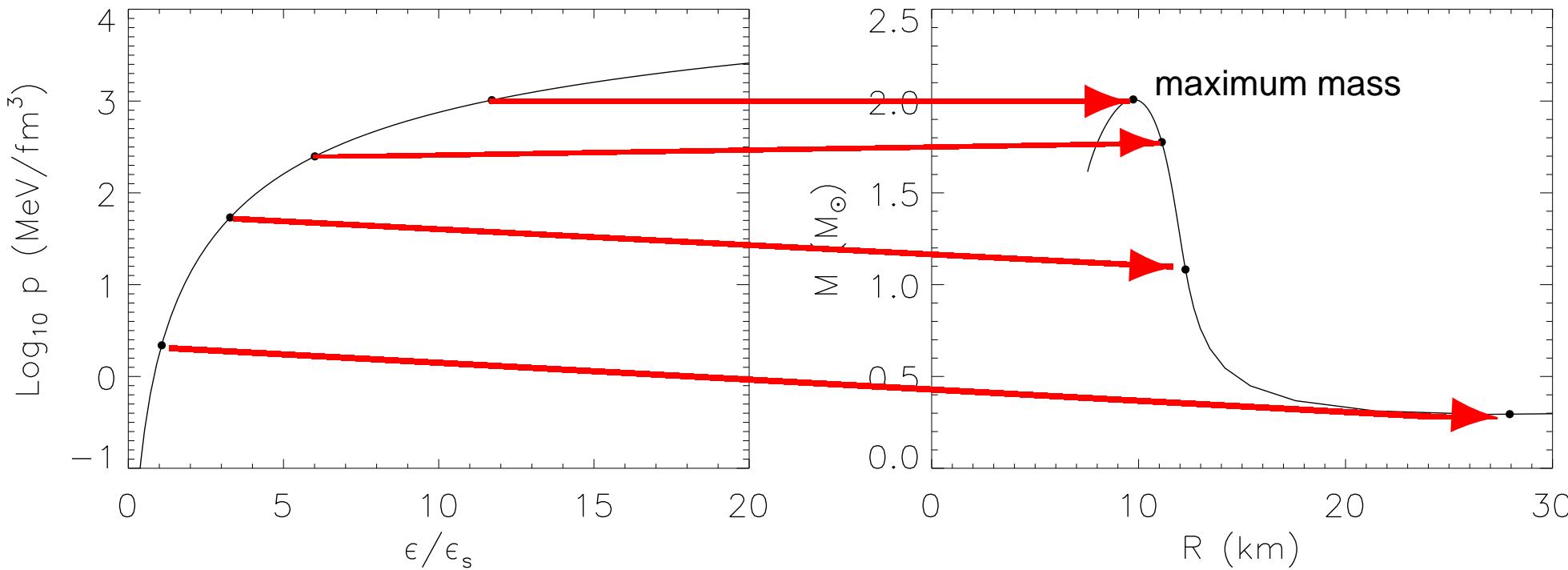
A NEUTRON STAR: SURFACE and INTERIOR



Neutron Star Structure

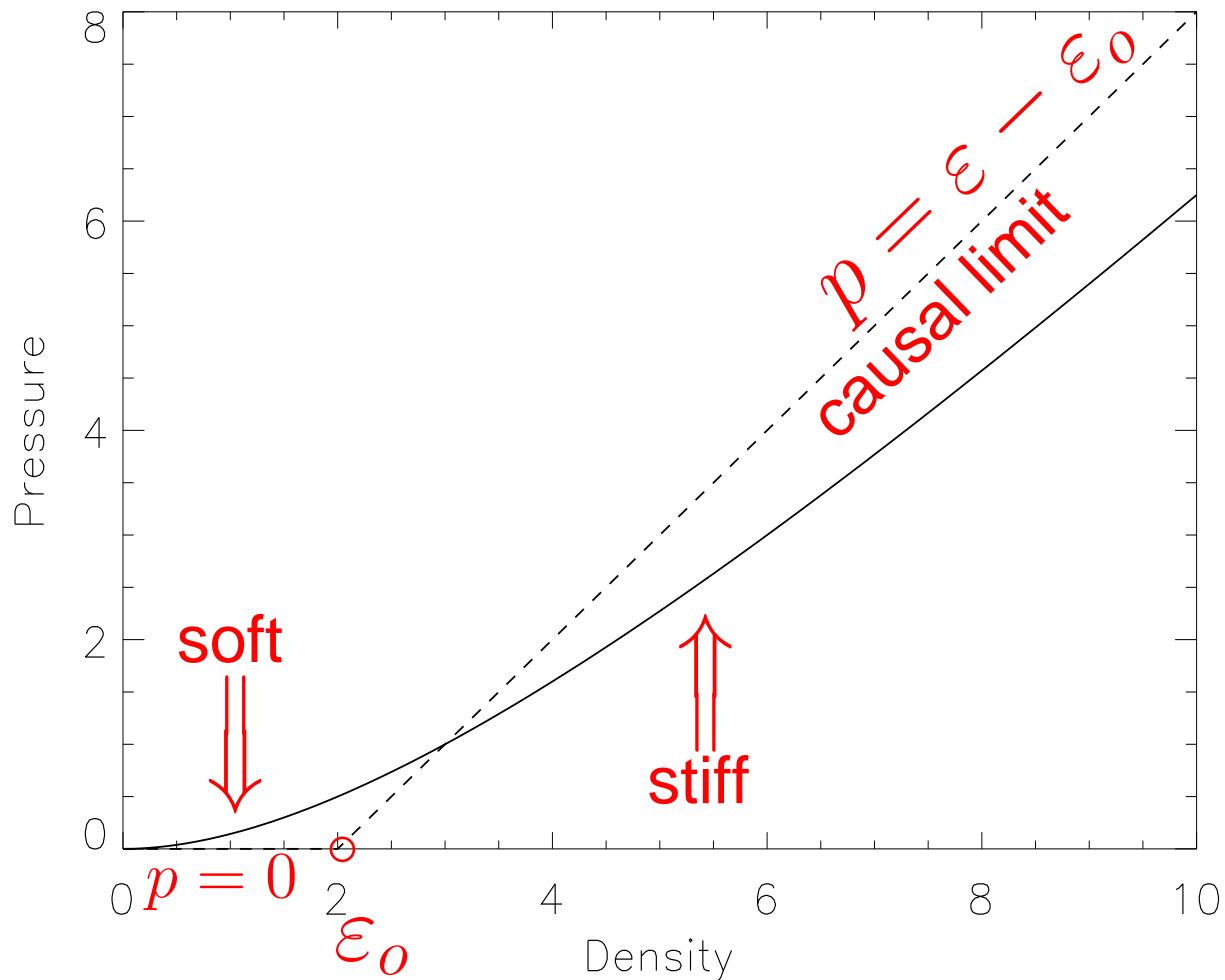
Tolman-Oppenheimer-Volkov equations

$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi pr^3)(\varepsilon + p)}{r(r - 2Gm/c^2)}$$
$$\frac{dm}{dr} = 4\pi \frac{\varepsilon}{c^2} r^2$$



Extreme Properties of Neutron Stars

The most compact configurations occur when the low-density equation of state is "soft" and the high-density equation of state is "stiff".



ϵ_0 is the only
EOS parameter

TOV solutions
scale with ϵ_0
even in 2-D

Koranda, Stergioulas
& Friedmann (1997)

Maximum Mass, Minimum Period

Theoretical limits from GR and causality

- $M_{max} = 4.2(\epsilon_s/\epsilon_f)^{1/2} M_\odot$ Rhoades & Ruffini (1974), Hartle (1978)
- $R_{min} = 2.9GM/c^2 = 4.3(M/M_\odot) \text{ km}$

Lindblom (1984), Glendenning (1992), Koranda, Stergioulas & Friedman (1997)

- $\epsilon_{central} < 4.5 \times 10^{15}(M_\odot/M_{largest})^2 \text{ g cm}^{-3}$
Lattimer & Prakash (2005)
- $P_{min} \simeq 0.74(M_\odot/M_{sph})^{1/2}(R_{sph}/10 \text{ km})^{3/2} \text{ ms}$
Koranda, Stergioulas & Friedman (1997)

- $P_{min} \simeq 0.96 \pm 0.03(M_\odot/M_{sph})^{1/2}(R_{sph}/10 \text{ km})^{3/2} \text{ ms}$
(empirical) Lattimer & Prakash (2004)
- $\epsilon_{central} > 0.91 \times 10^{15}(1 \text{ ms}/P_{min})^2 \text{ g cm}^{-3}$ (empirical)
- $cJ/GM^2 \lesssim 0.5$ (empirical, neutron star)

Mass-Radius Diagram and Theoretical Constraints

GR:

$$R > 2GM/c^2$$

$P < \infty :$

$$R > (9/4)GM/c^2$$

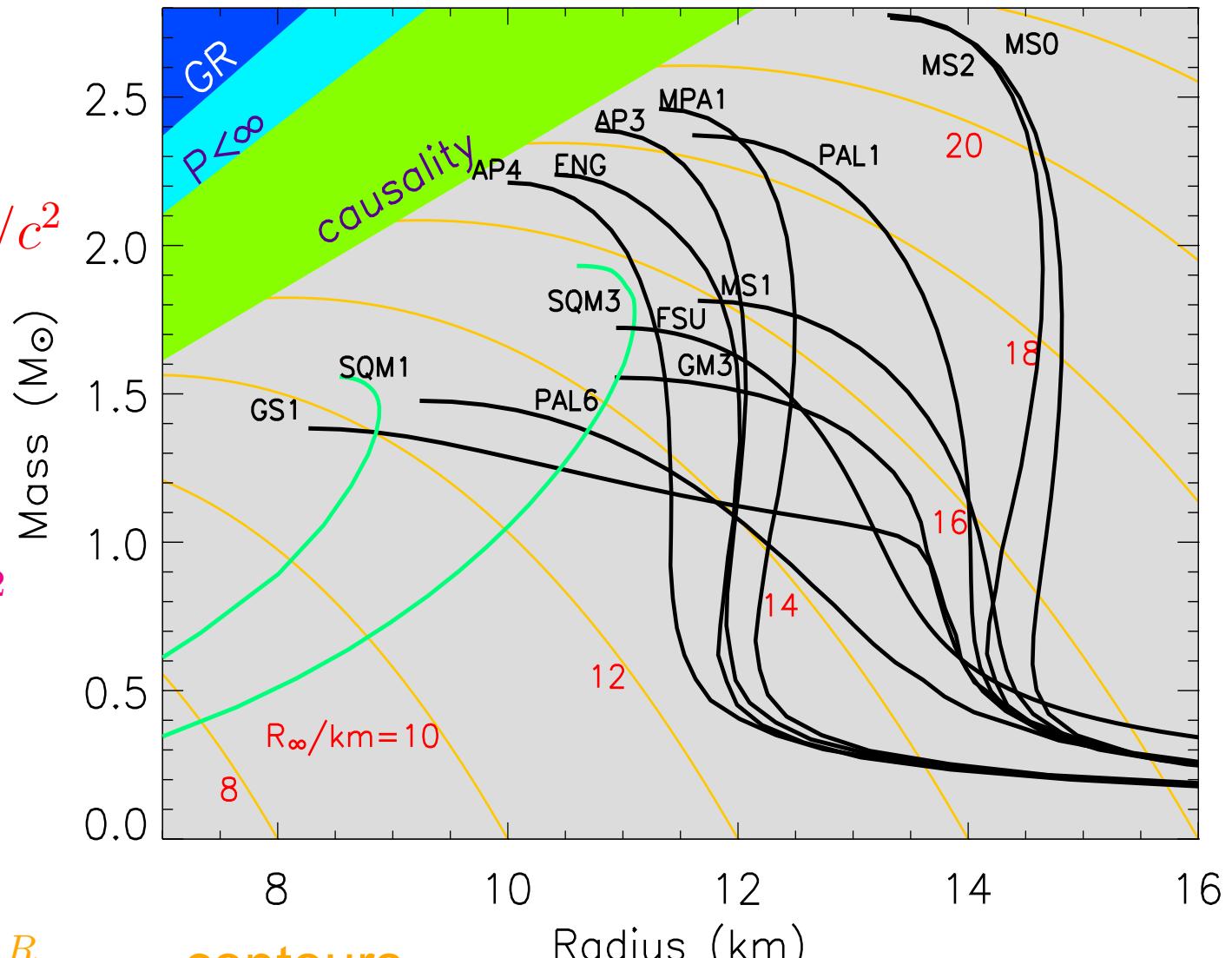
causality:

$$R \gtrsim 2.9GM/c^2$$

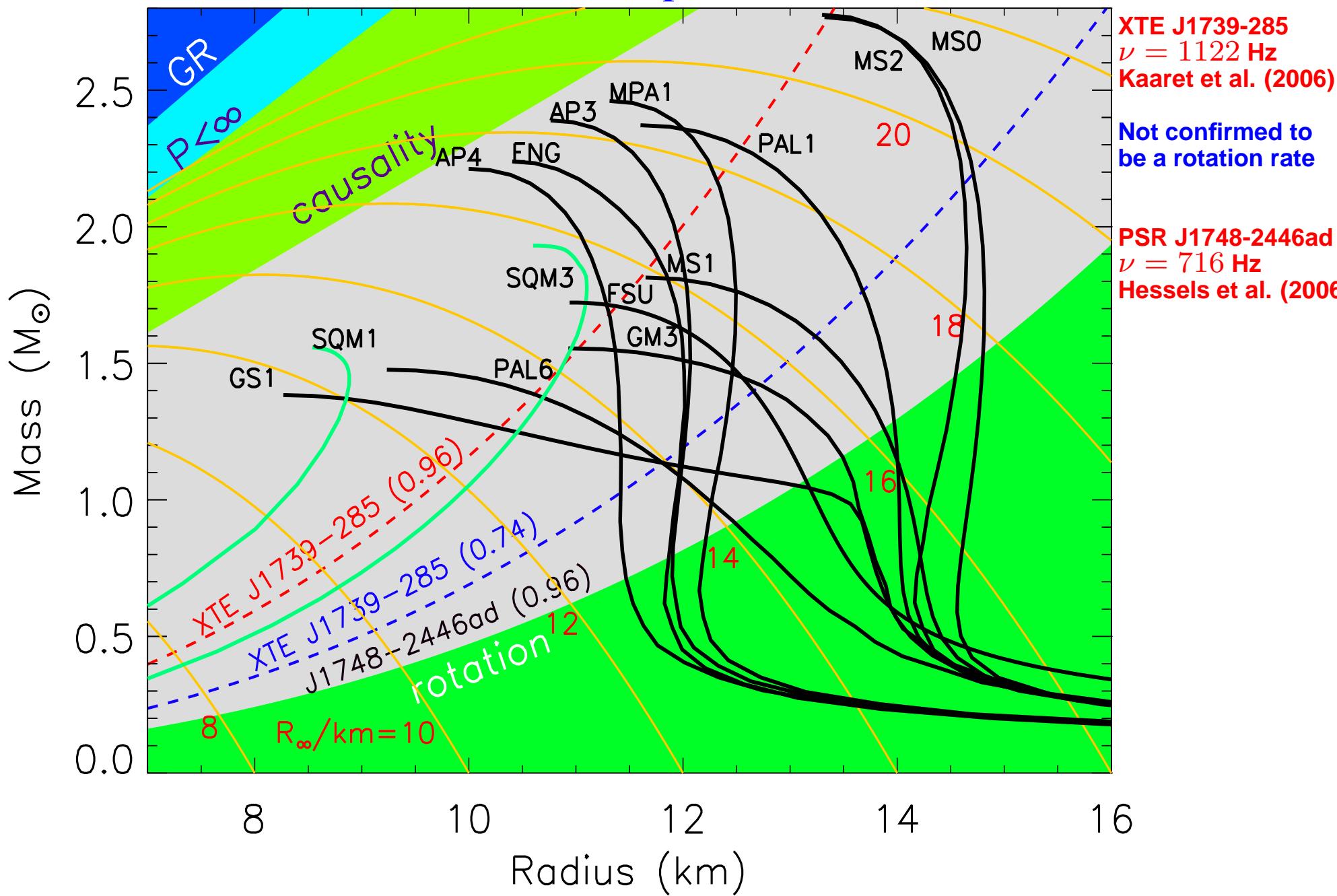
— normal NS

— SQS

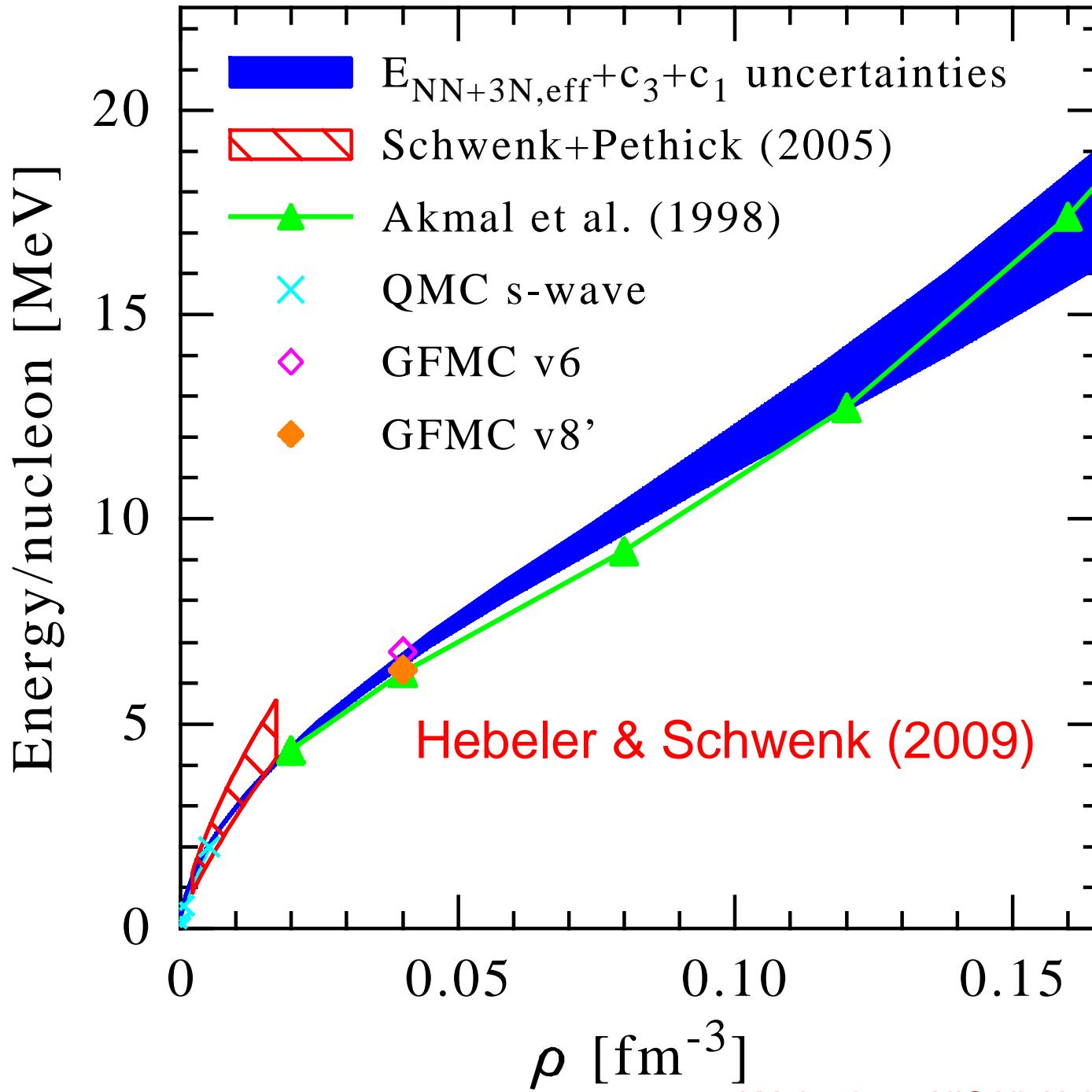
— $R_\infty = \frac{R}{\sqrt{1-2GM/Rc^2}}$ contours



Constraints from Pulsar Spins



Pure Neutron Matter



Estimating Symmetry Parameters From Neutron Matter Results

$$E_n = E(n_s, x = 0) \simeq 16.3 \pm 2.1 \text{ MeV}, \quad P_n \simeq 2.5 \pm 0.7 \text{ MeV fm}^{-3}$$

$$S'_v \equiv \left. \frac{\partial \ln E_{sym}(n)}{\partial \ln n} \right|_{n_s} = P_n/n_s = 15.6 \pm 4.4 \text{ MeV}$$

- Simple Model

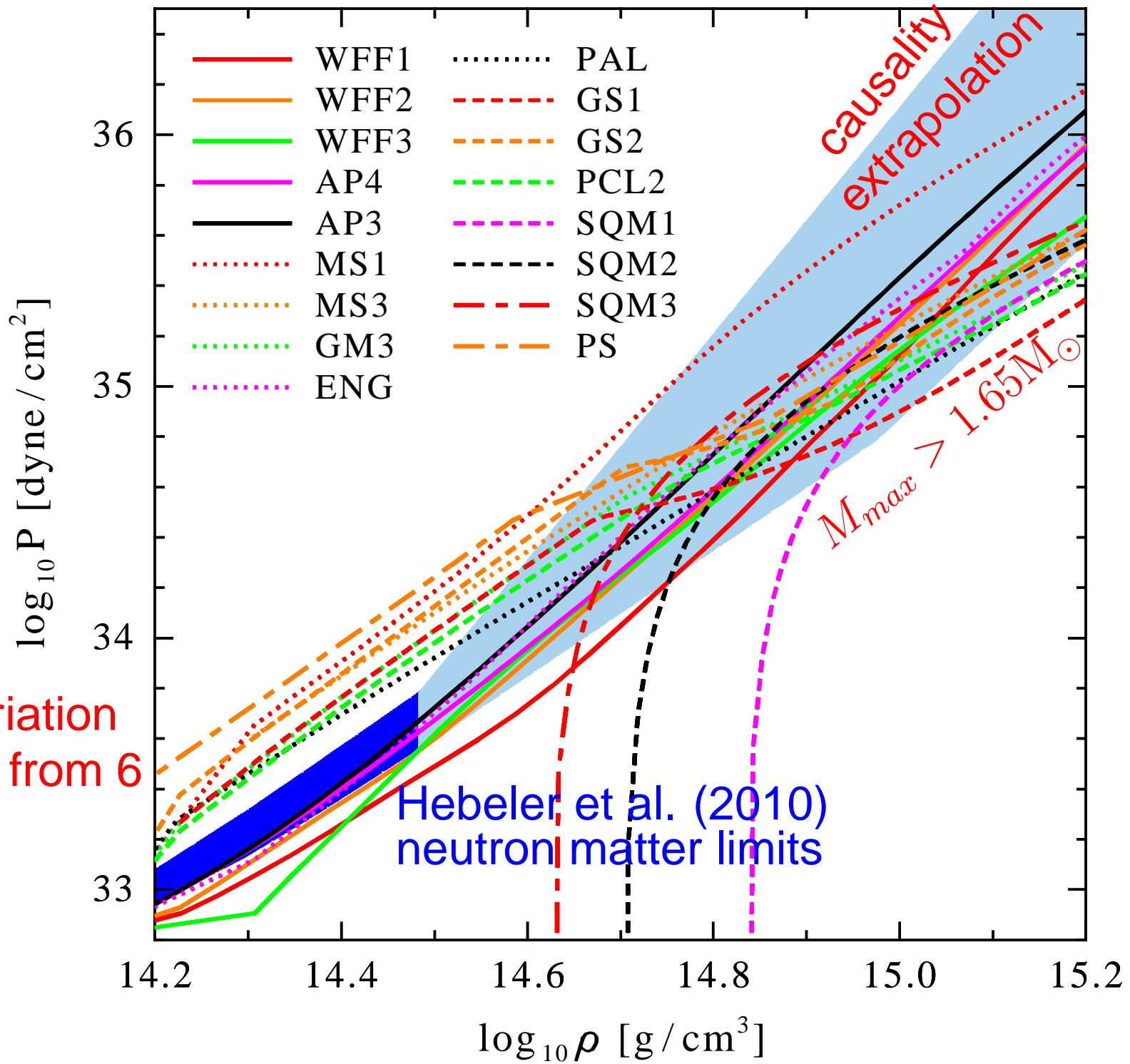
$$E_{sym}(n) = S_v (n/n_s)^p$$

$$S_v = E_n + B \simeq 32.3 \pm 2.1 \text{ MeV}, \quad p = S'_v/S_v \simeq 0.48 \pm 0.14$$

- More Accurate Model

$$E_{sym}(n) = S_k (n/n_s)^{2/3} + (S_v - S_k) (n/n_s)^\gamma$$

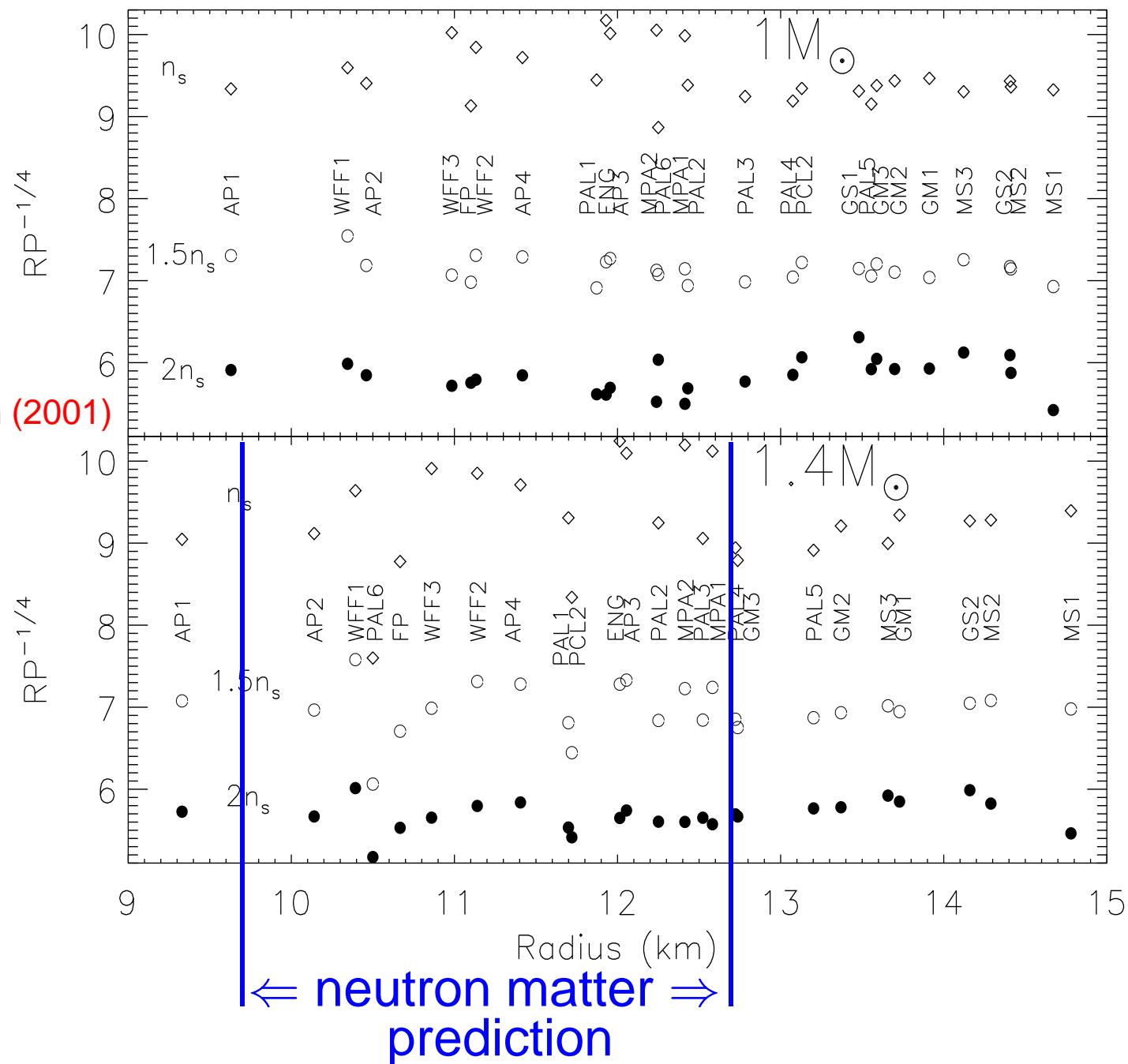
$$S_k \simeq 17 \text{ MeV}, \quad \gamma = \frac{S'_v - 2S_v/3}{S_v - S_k} \simeq 0.28 \pm 0.29$$



The Radius – Pressure Correlation

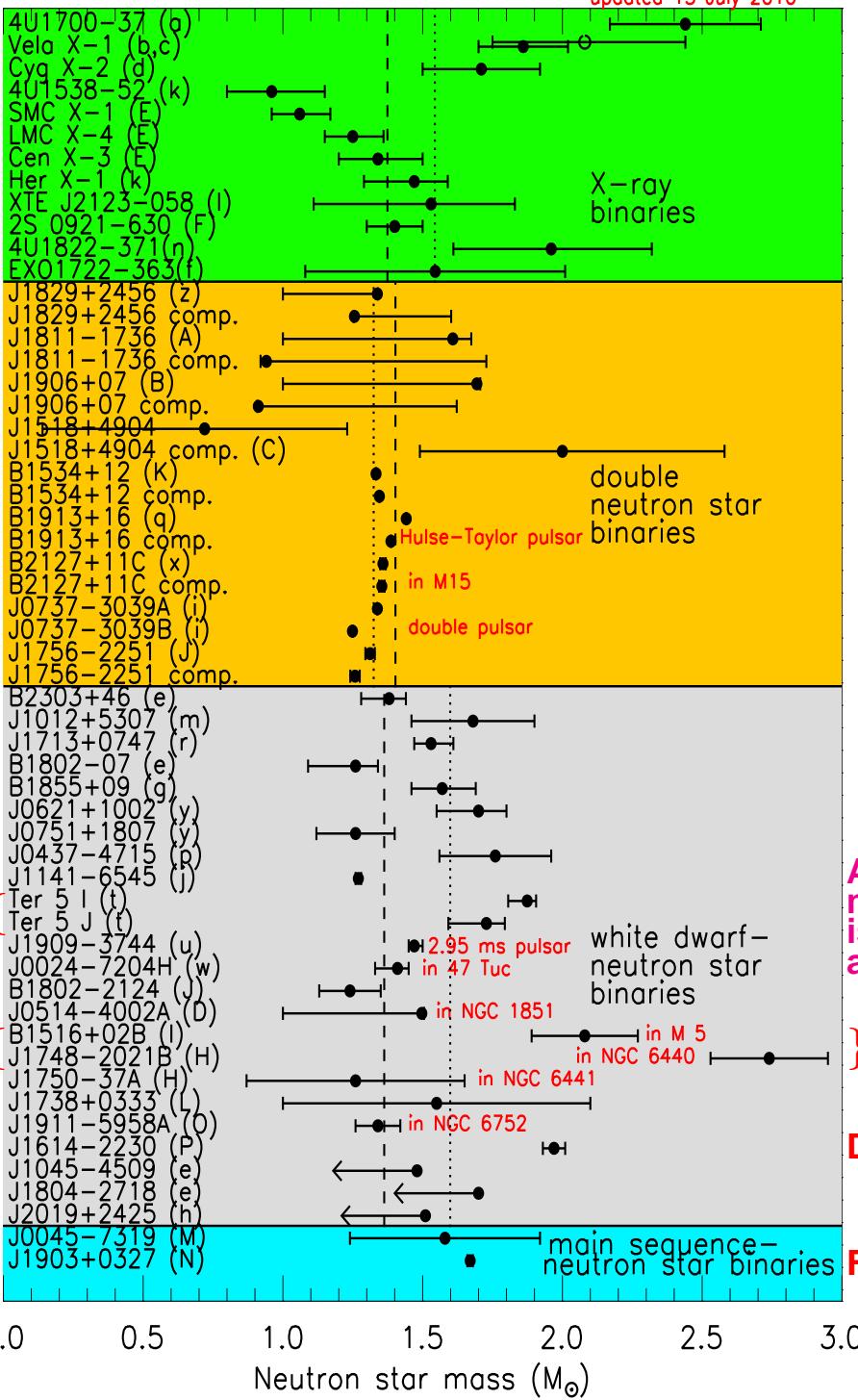
$$R \propto p^{1/4}$$

Lattimer & Prakash (2001)



Black hole? \Rightarrow

updated 13 July 2010



Possible Kinds of Observations

- Maximum and Minimum Mass (binary pulsars)
- Minimum Rotational Period*
- Radiation Radii or Redshifts from X-ray Thermal Emission*
- Crustal Cooling Timescale from X-ray Transients*
- X-ray Bursts from Accreting Neutron Stars*
- Seismology from Giant Flares in SGR's*
- Neutron Star Thermal Evolution (URCA or not)*
- Moments of Inertia from Spin-Orbit Coupling*
- Neutrinos from Proto-Neutron Stars (Binding Energies, Neutrino Opacities, Radii)*
- Pulse Shape Modulations*
- Gravitational Radiation from Neutron Star Mergers*
(Masses, Radii from tidal Love numbers)

* Significant dependence on symmetry energy

Potentially Observable Quantities

- Apparent angular diameter from flux and temperature measurements

$$\beta \equiv GM/Rc^2$$

$$\frac{R_\infty}{D} = \frac{R}{D} \frac{1}{\sqrt{1 - 2\beta}} = \sqrt{\frac{F_\infty}{\sigma}} \frac{1}{f_c^2 T_\infty^2}$$

- Redshift

$$z = (1 - 2\beta)^{-1/2} - 1$$

- Eddington flux

$$F_{Edd} = \frac{GMc}{\kappa c^2 D^2} (1 - 2\beta)^{1/2}$$

- Crust thickness

$$\frac{m_b c^2}{2} \ln \mathcal{H} \equiv \int_0^{p_t} \frac{dp}{n} = \mu_{n,t} - \mu_{n,t}(p=0)$$

$$\frac{\Delta}{R} \equiv \frac{R - R_t}{R} \simeq \frac{(\mathcal{H} - 1)(1 - 2\beta)}{\mathcal{H} + 2\beta - 1} \simeq (\mathcal{H} - 1) \left(\frac{1}{2\beta} - 1 \right).$$

- Moment of Inertia

$$I \simeq (0.237 \pm 0.008) MR^2 (1 + 2.84\beta + 18.9\beta^4) M_\odot \text{ km}^2$$

- Crustal Moment of Inertia

$$\frac{\Delta I}{I} \simeq \frac{8\pi}{3} \frac{R^6 p_t}{IMc^2}$$

- Binding Energy

$$\frac{\text{B.E.}}{Mc^2} \simeq (0.60 \pm 0.05) \frac{\beta}{1 - \beta/2}$$

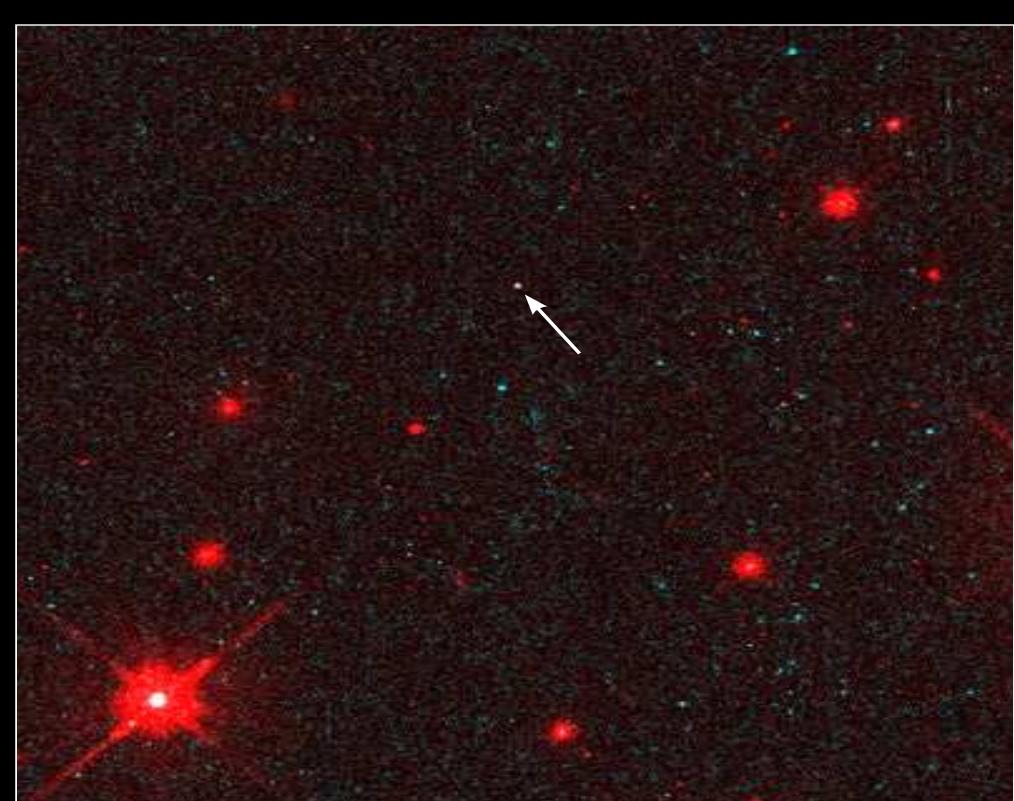
Radiation Radius

- Combination of flux and temperature measurements yields apparent angular diameter (pseudo-BB):

$$\frac{R_\infty}{d} = \frac{R}{d} \frac{1}{\sqrt{1 - 2GM/Rc^2}}$$

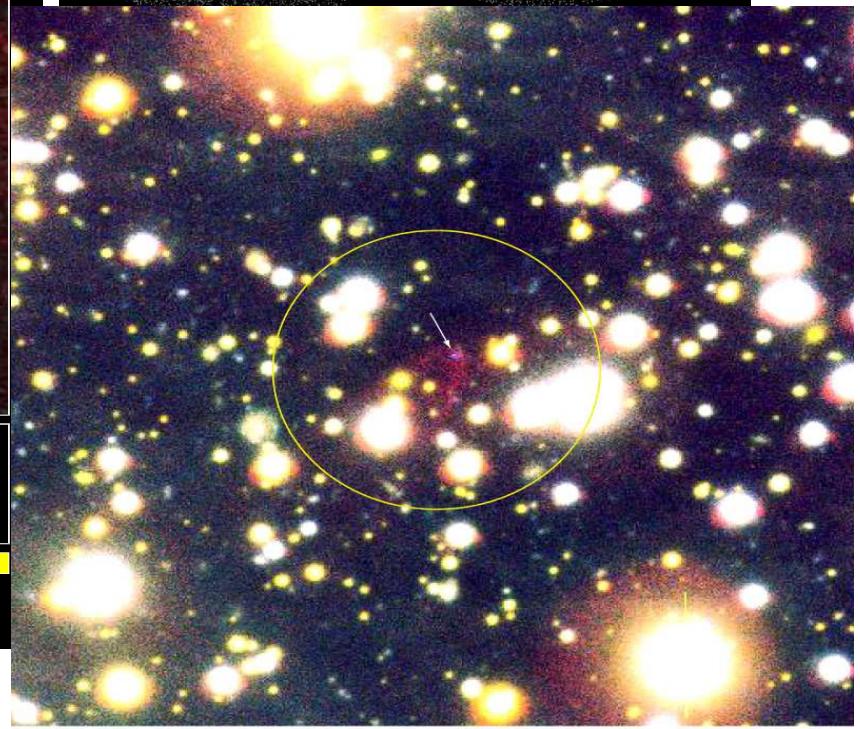
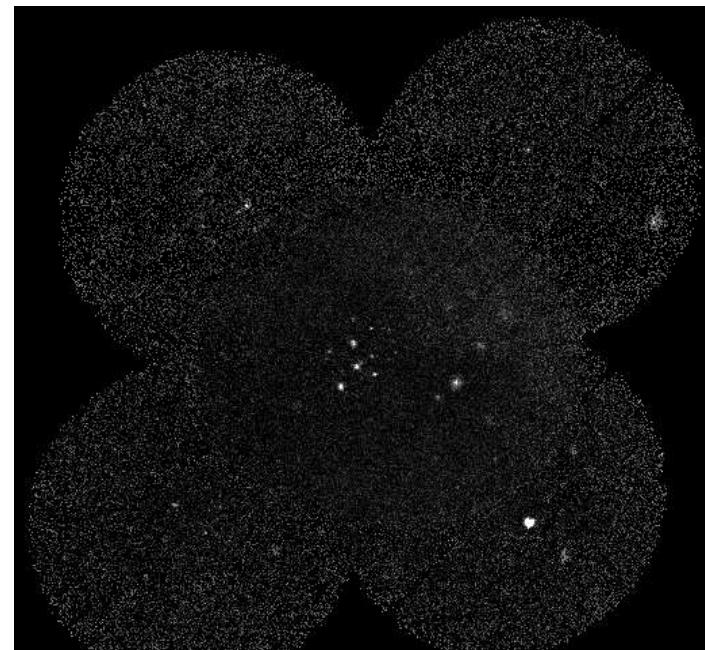
- Observational uncertainties include distance, interstellar H absorption (hard UV and X-rays), atmospheric composition
- Best chances for accurate radii are from
 - Nearby isolated neutron stars (parallax measurable)
 - Quiescent X-ray binaries in globular clusters (reliable distances, low B H-atmospheres)

RX J1856-3754



Isolated Neutron Star RX J185635-3754
Hubble Space Telescope • WFPC2

PRC97-32 • ST Scl OPO • September 25, 1997
F. Walter (State University of New York at Stony Brook) and NASA



A Bowshock Nebula Near the Neutron Star RX J1856.5-3754 (Detail)
(VLT KUEYEN + FORS2)

ESO PR Photo 23b/00 (11 September 2000)

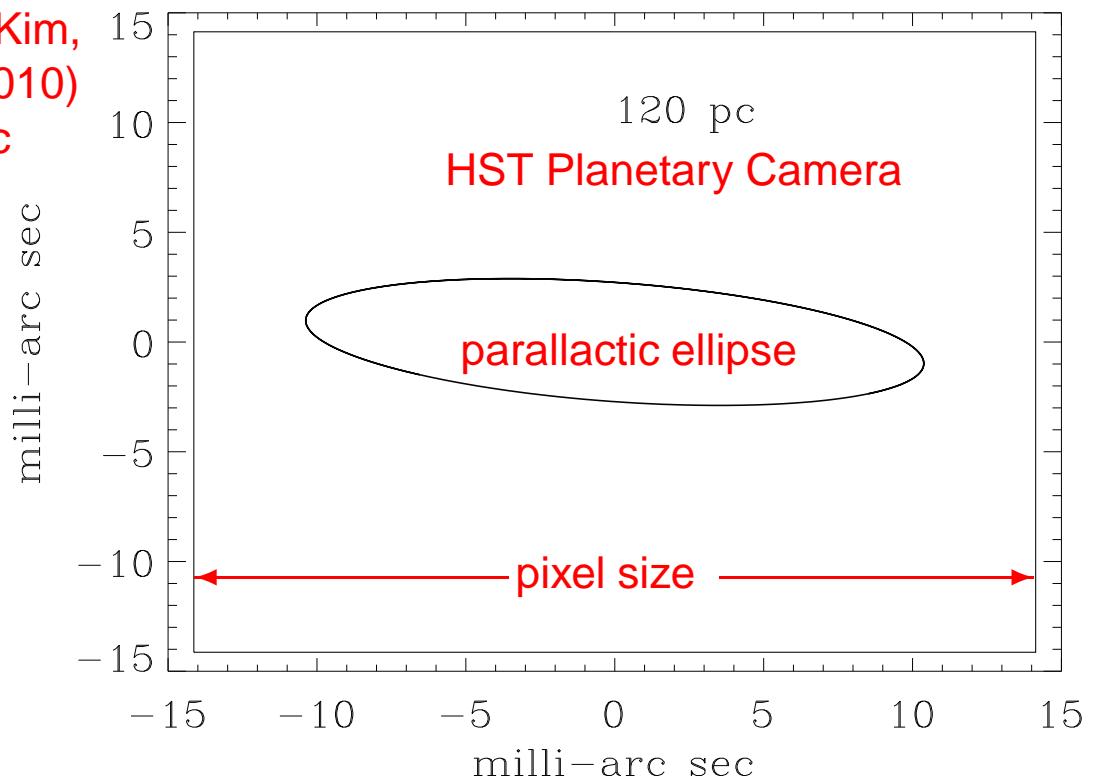
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J.M. Lattimer, NIC XI, Heidelberg, 20/07/10 – p. 17/31



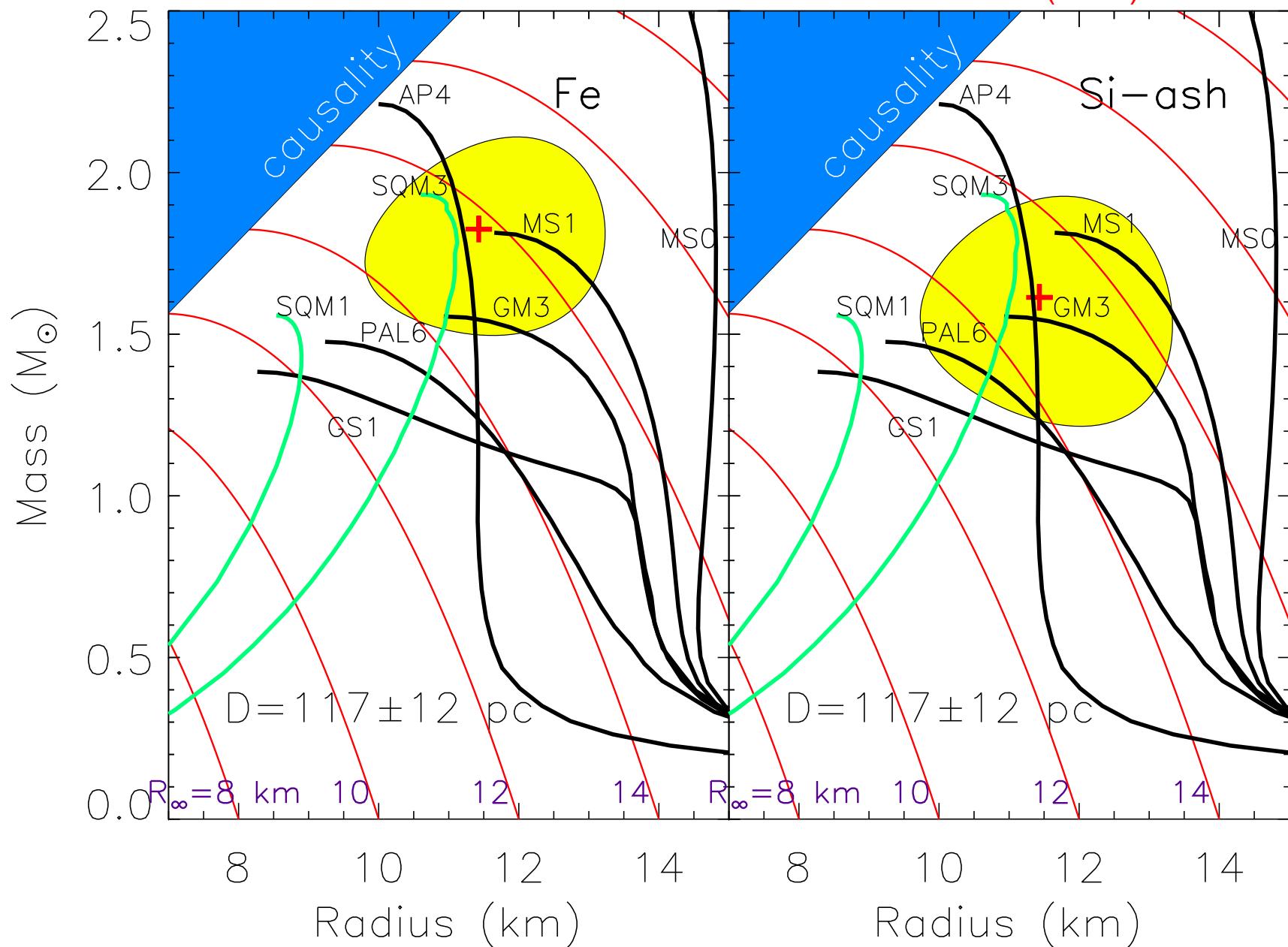
Astrometry of RXJ 1856-3754

- Walter & Lattimer (2002) determined $D = 117 \pm 12$ pc and $v \simeq 190$ km/s from 1996-1999 HST Planetary Camera observations
- Star's age is probably 0.5 million years, origin Cor Aus
- Kaplan, van Kerkwijk & Anderson (2002): $D = 140 \pm 40$ pc using same data
- van Kerkwijk & Kaplan (2007, conference proceeding) revised this to $D = 161 \pm 16$ pc based on 2002-2004 High-Resolution Camera of the Advanced Camera for Surveys HST observations (double the resolution)
- Walter, Eisenbeiß, Lattimer, Kim, Hambaryan & Neuhäuser (2010) determined $D \simeq 122 \pm 13$ pc with 2002-2004 HST data

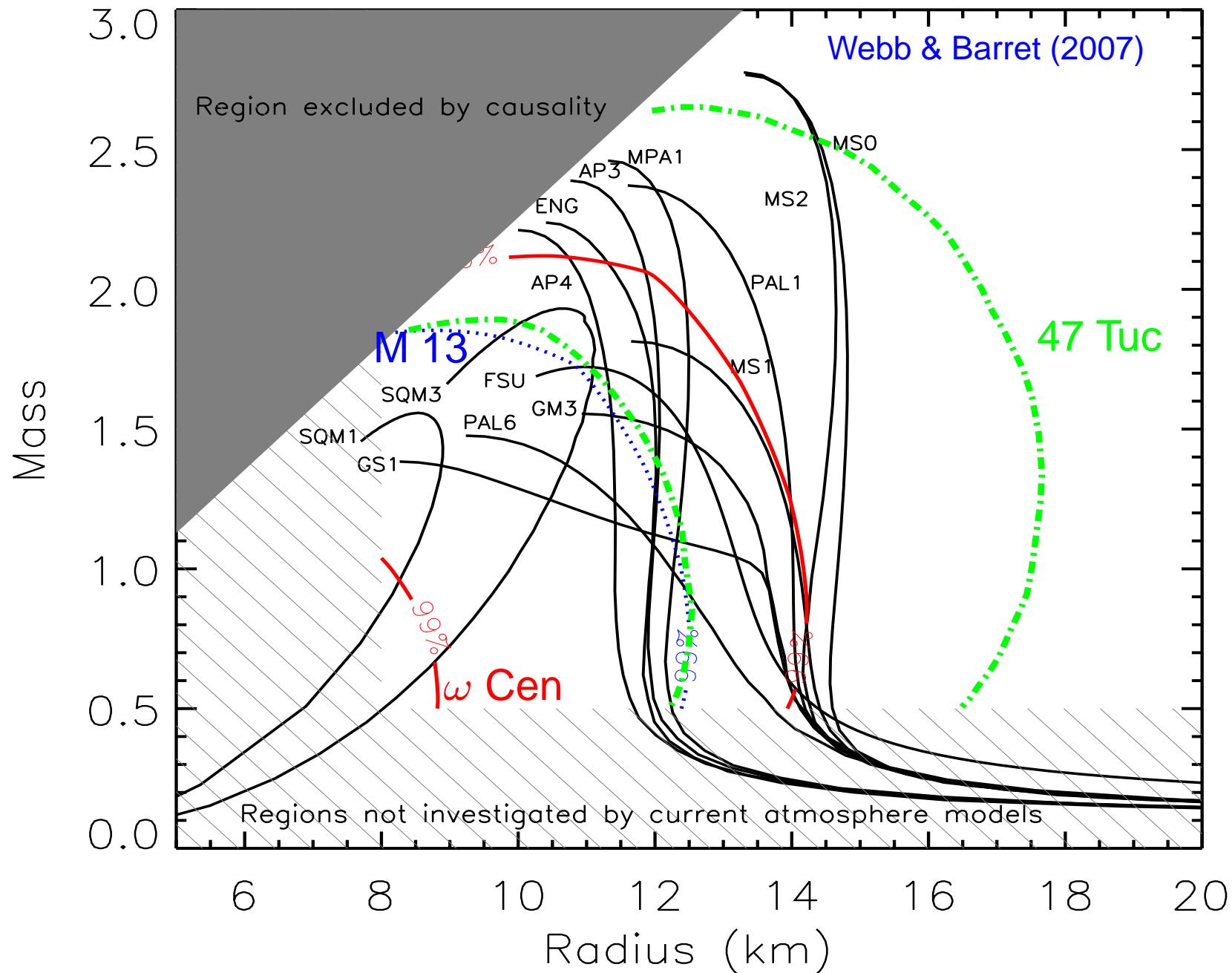


RX J1856-3754

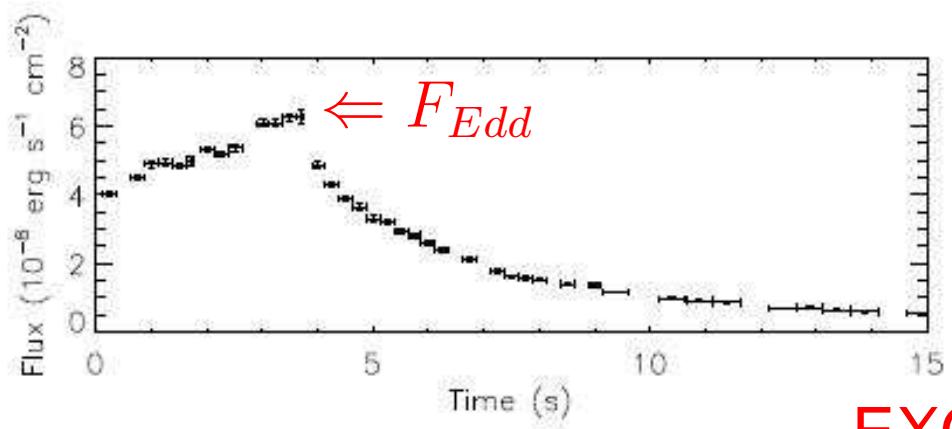
Walter & Lattimer (2002)



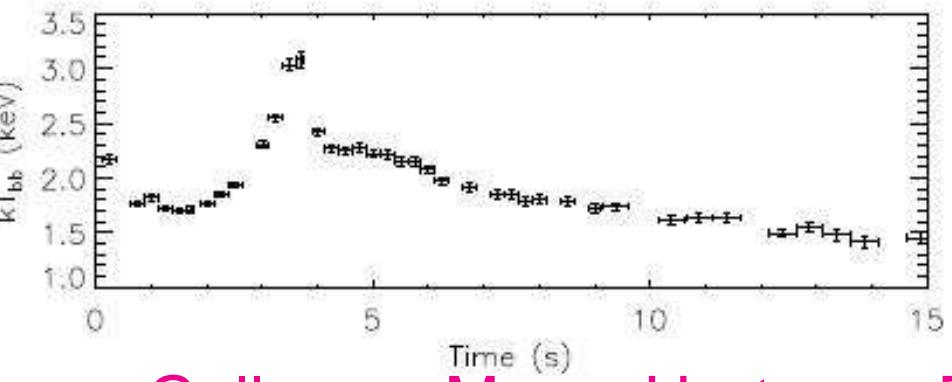
Radiation Radius: Globular Cluster Sources



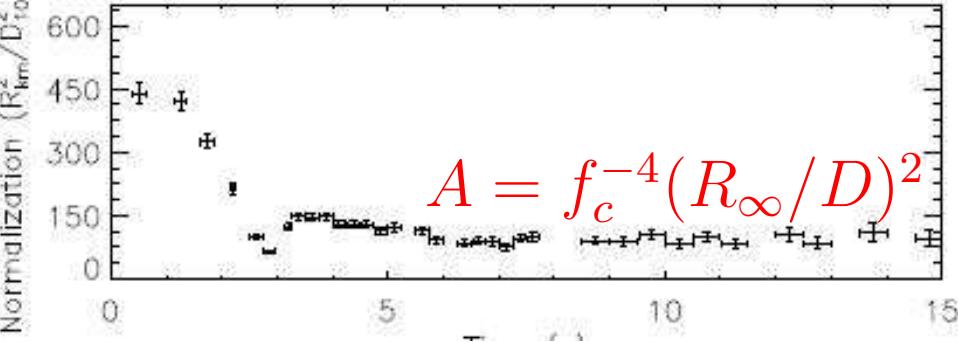
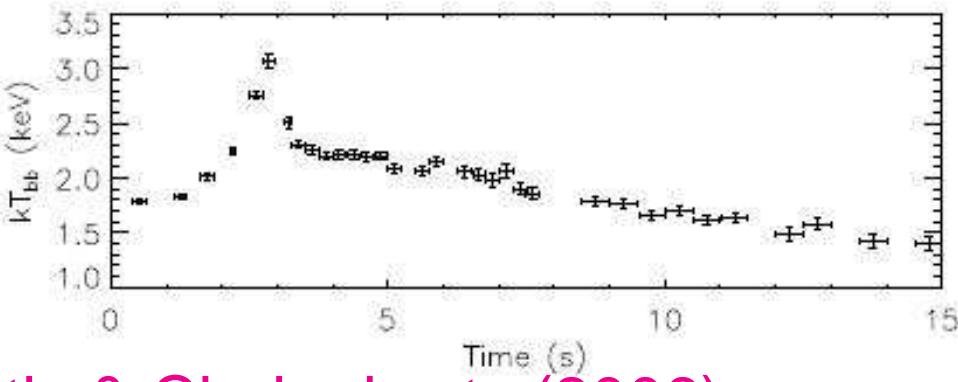
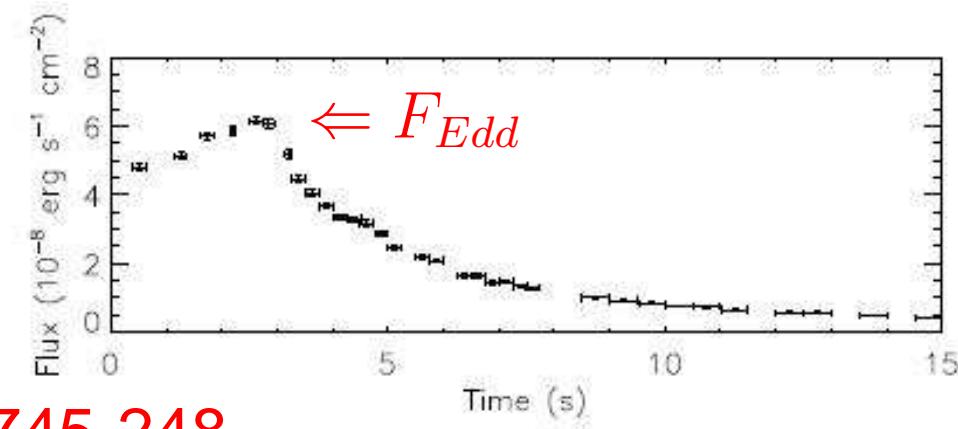
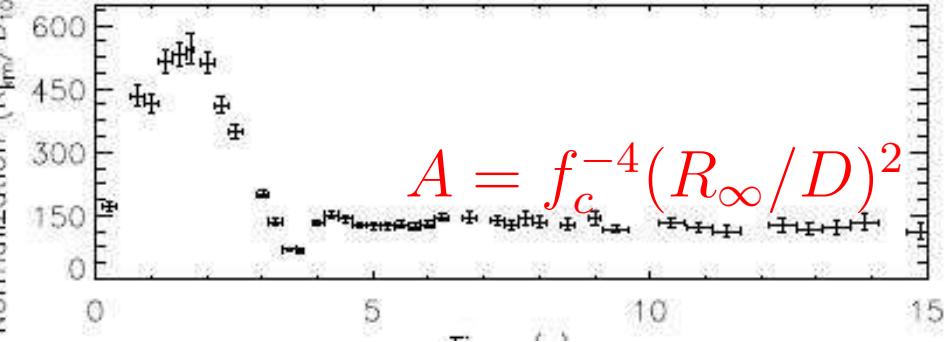
Photospheric Radius Expansion X-Ray Bursts



EXO 1745-248



Galloway, Muno, Hartman, Psaltis & Chakrabarty (2006)



Systematics assuming $R_{ph} = R$

$$F_{Edd} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\frac{GM}{R_{ph}c^2}} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\beta}$$

$$\kappa \simeq 0.2(1 + X) \text{ cm}^2\text{g}^{-1}$$

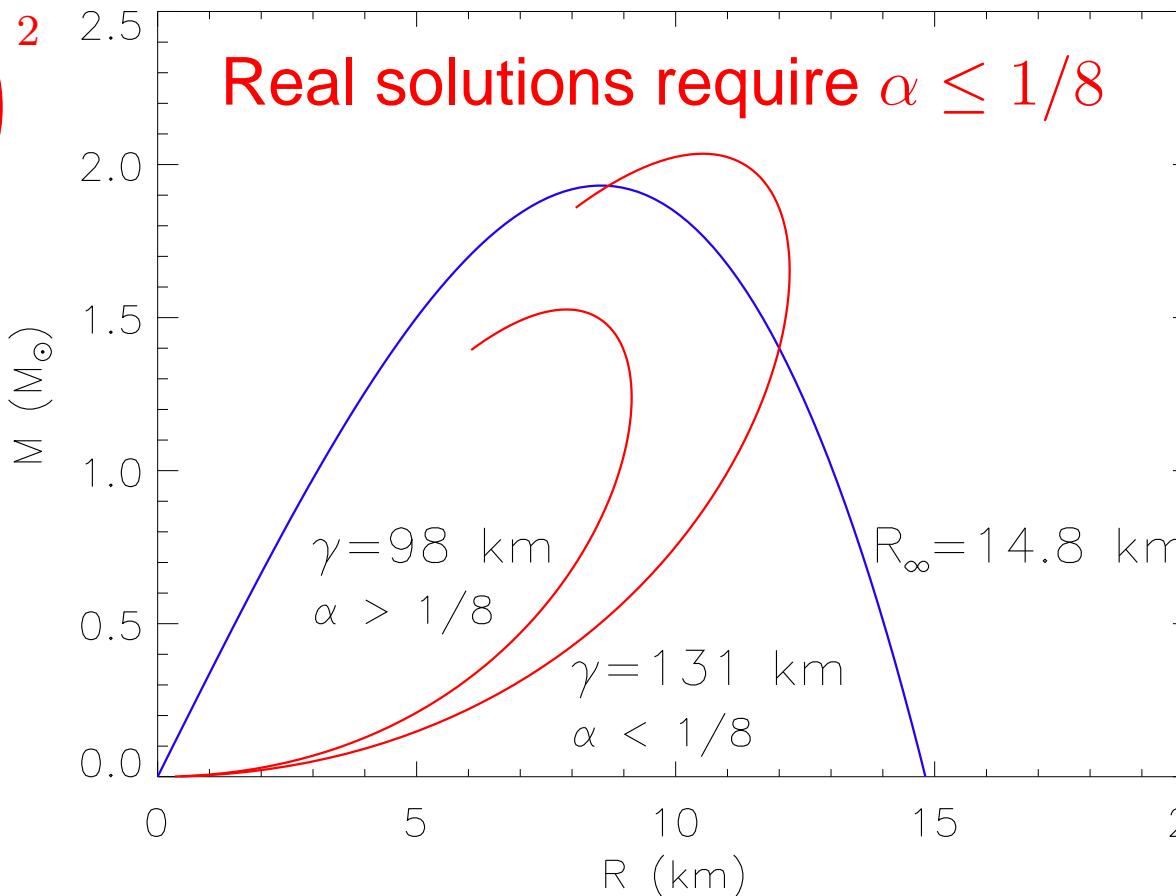
$$A = \frac{F_\infty}{\sigma T_\infty^4} = f_c^{-4} \left(\frac{R_\infty}{D} \right)^2$$

$$\alpha = \frac{F_{Edd}}{\sqrt{A}} \frac{\kappa D}{c^3 f_c^2} = \beta(1 - 2\beta)$$

$$\gamma = \frac{Ac^3 f_c^4}{F_{Edd}\kappa} = \frac{R}{\beta(1 - 2\beta)^{3/2}}$$

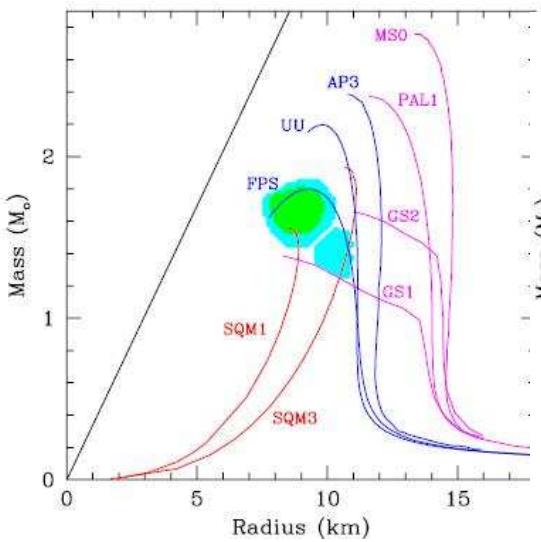
$$\beta = \frac{1}{4} \pm \frac{1}{4}\sqrt{1 - 8\alpha}$$

$$R_\infty = \alpha\gamma$$

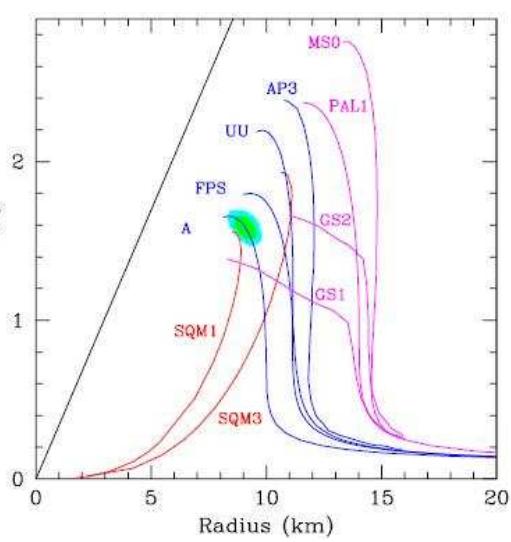


$$R_{ph} = R$$

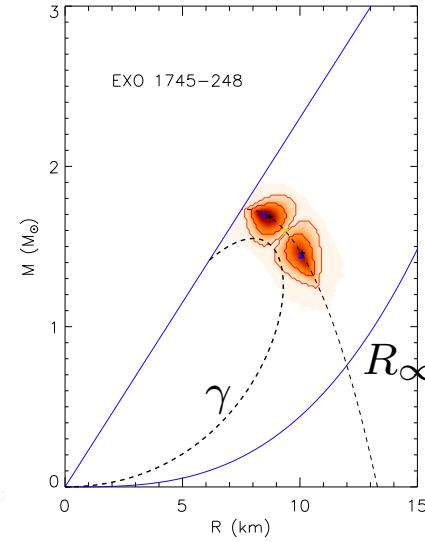
EXO 1745-248
Özel et al. (2009)



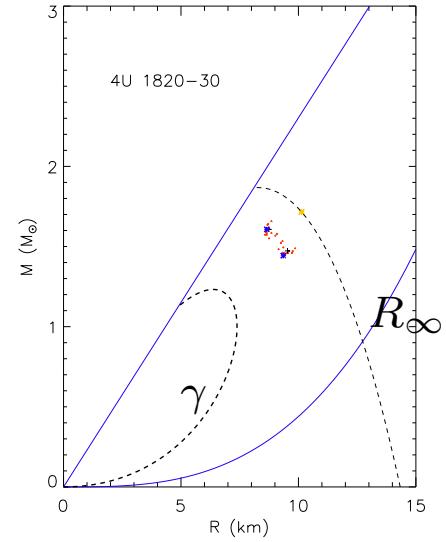
4U 1820-30
Güver et al. (2010)



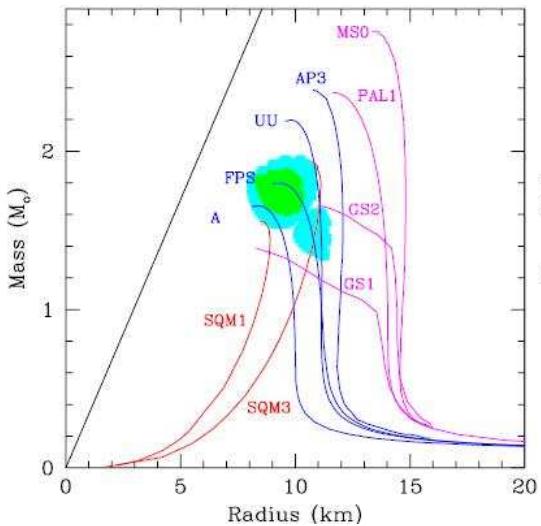
EXO 1745-248
 $\alpha = 0.14 \pm 0.02$



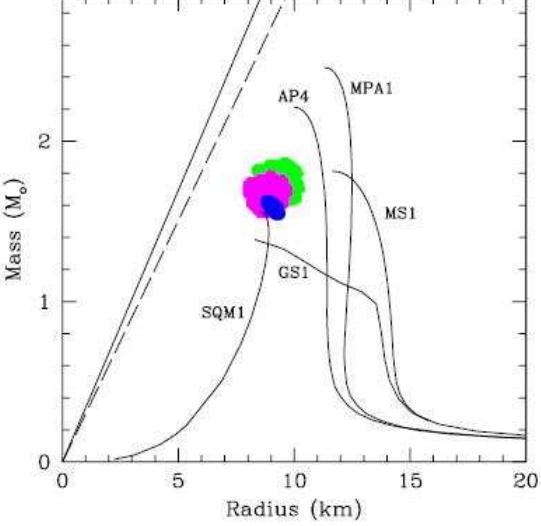
4U 1820-30
 $\alpha = 0.18 \pm 0.02$



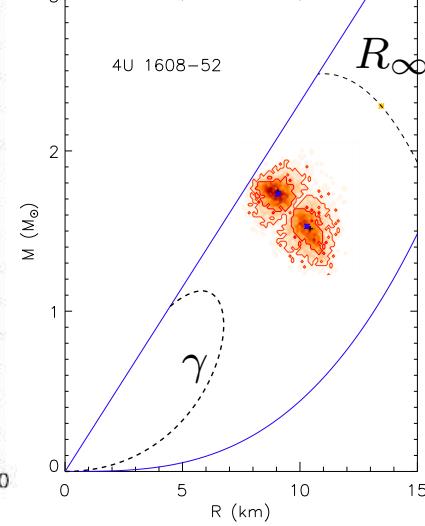
4U 1608-52
Güver et al. (2010)



Özel, Baym & Güver (2010)



4U 1608-52
 $\alpha = 0.26 \pm 0.11$



Systematics assuming $R_{ph} \gg R$

$$F_{Edd} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\frac{GM}{R_{ph}c^2}} \simeq \frac{GMc}{\kappa D^2}$$

$$\kappa \simeq 0.2(1 + X) \text{ cm}^2\text{g}^{-1}$$

$$A = \frac{F_\infty}{\sigma T_\infty^4} = f_c^{-4} \left(\frac{R_\infty}{D} \right)^2$$

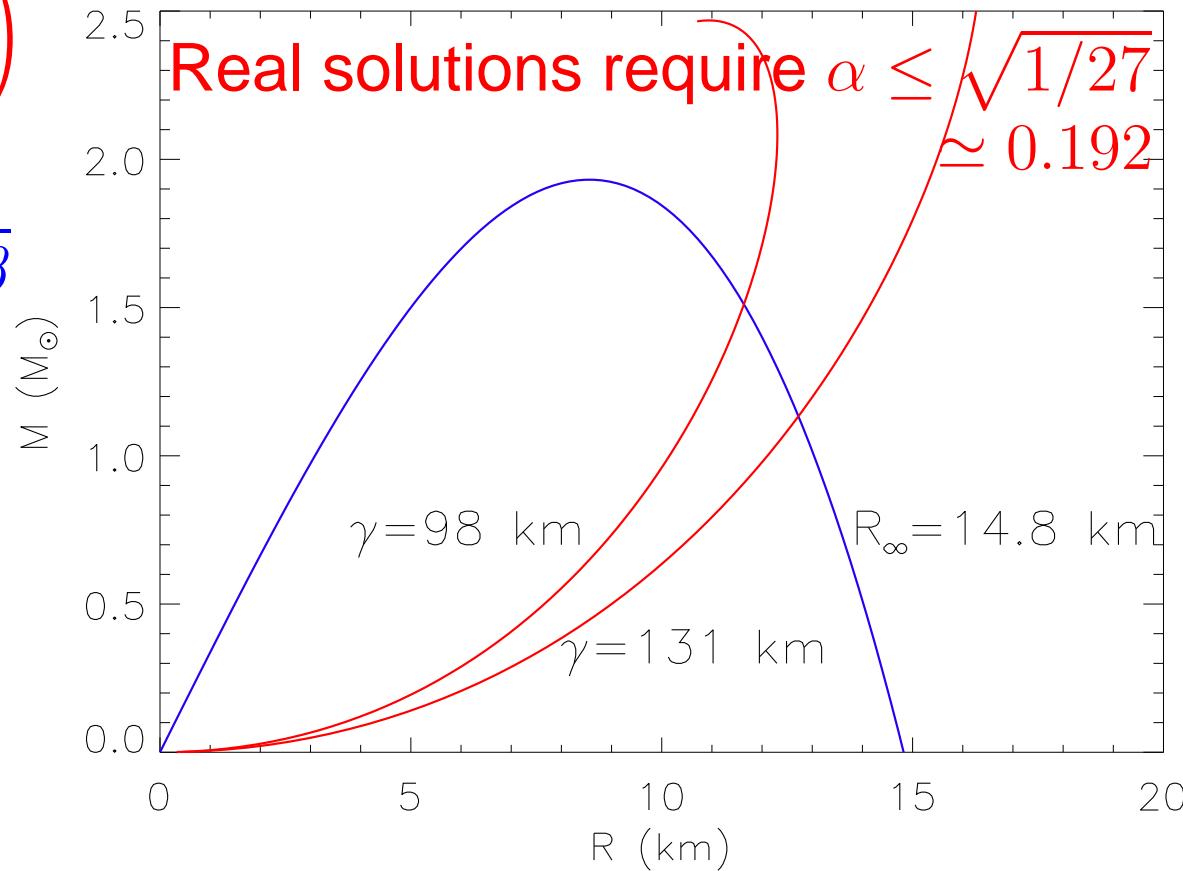
$$\alpha = \frac{F_{Edd}}{\sqrt{A}} \frac{\kappa D}{c^3 f_c^2} = \beta \sqrt{1 - 2\beta}$$

$$\gamma = \frac{Ac^3 f_c^4}{F_{Edd} \kappa} = \frac{R}{\beta(1 - 2\beta)}$$

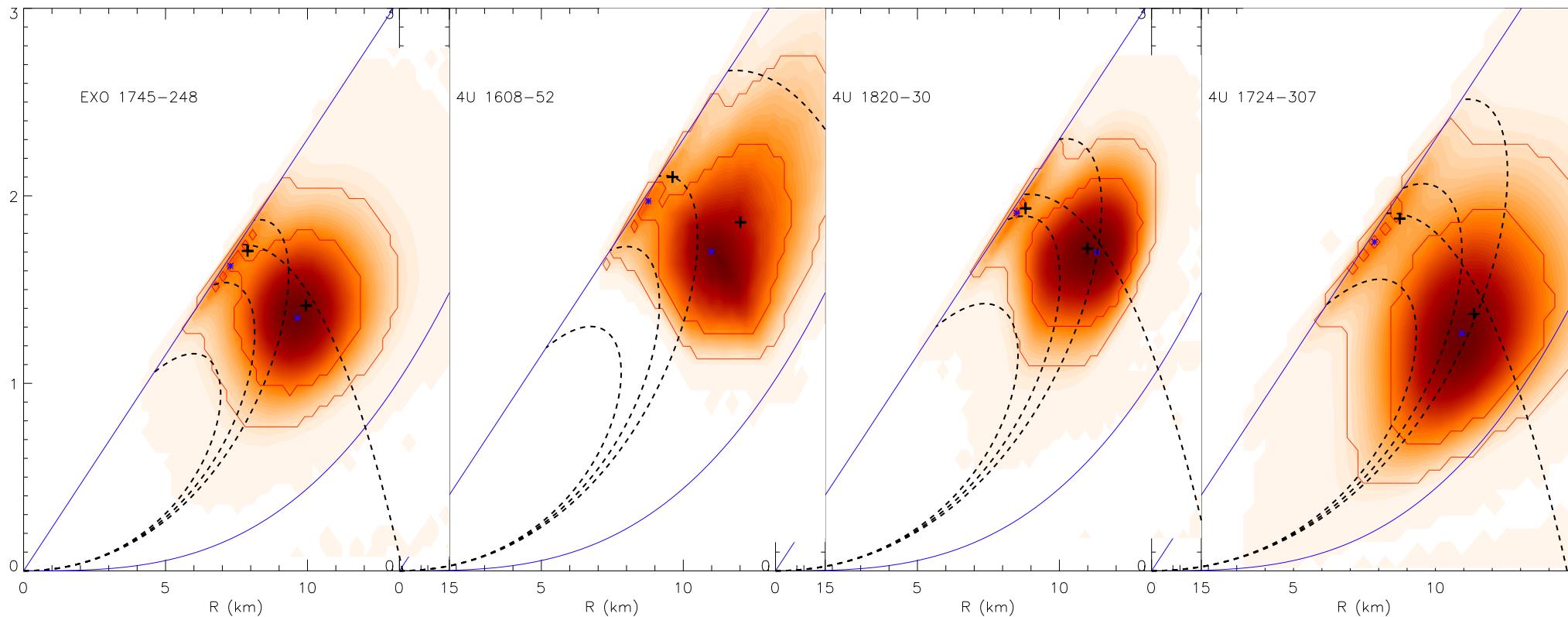
$$\beta = \frac{1}{6} [1 + \sqrt{3} \sin \theta - \cos \theta]$$

$$\theta = \frac{1}{3} \cos^{-1}(1 - 54\alpha^2)$$

$$R_\infty = \alpha \gamma$$



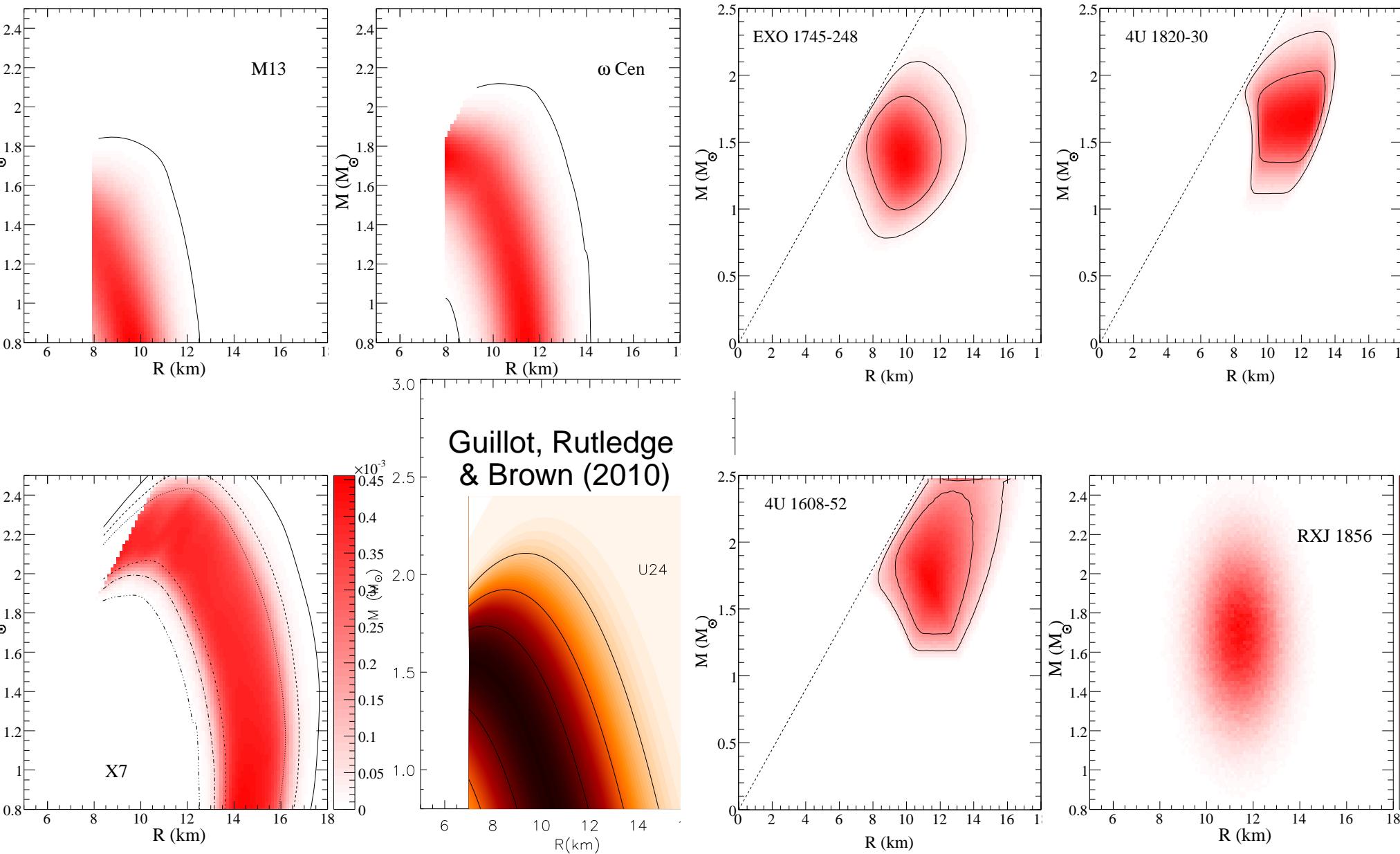
Arbitrary r_{ph}/R



Sources of observational data:

- 4U 1745-248: Özel et al. (2009)
- 4U 1608-52: Güver et al. (2010)
- 4U 1820-30: Güver et al. (2010)
- 4U 1724-307: Suleimanov et al. (2010)

M-R Probability Distributions



Bayesian Analysis

$P(\mathcal{M}_i|D)$: conditional probability of the model \mathcal{M}_i given the data D (what we want).

$$P(\mathcal{M}_i|D) = \frac{P(D|\mathcal{M}_i)P(\mathcal{M}_i)}{\sum_k P(D|\mathcal{M}_k)P(\mathcal{M}_k)} \propto P(D|\mathcal{M}_i).$$

$P(\mathcal{M}_i)$: prior probability of the model \mathcal{M}_i without any information from the data D (assumed uniform).

$P(D|\mathcal{M}_i)$: conditional probability of the data D given the model \mathcal{M}_i .
 i is the EOS parameter set and the set of specific masses j .

$$P(D|\mathcal{M}_i) \propto \prod_{k,j} \mathcal{D}_{k,j}|_{M=M_j, R=R_k(M_j)}$$

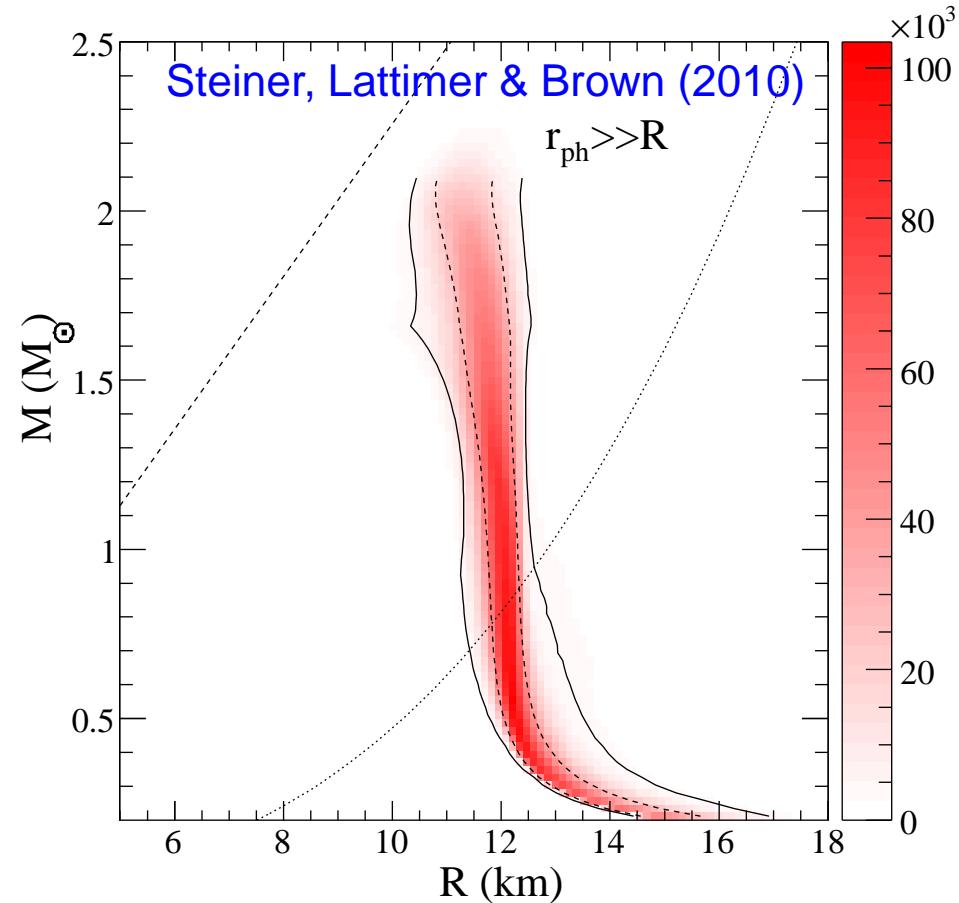
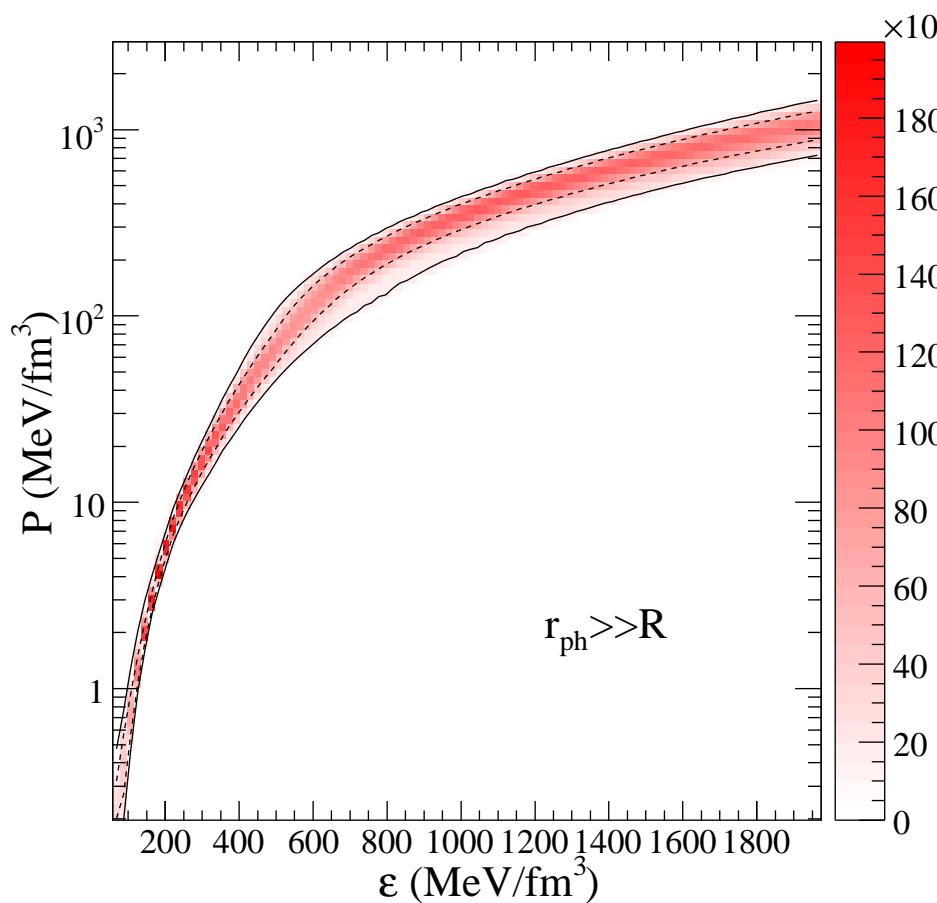
$\mathcal{D}_{k,j}$: probability of observing the radius R_j given a mass M_j in observation k ($M - R$ probabilities).

Posterior probability for a model parameter p_i (marginal estimation):

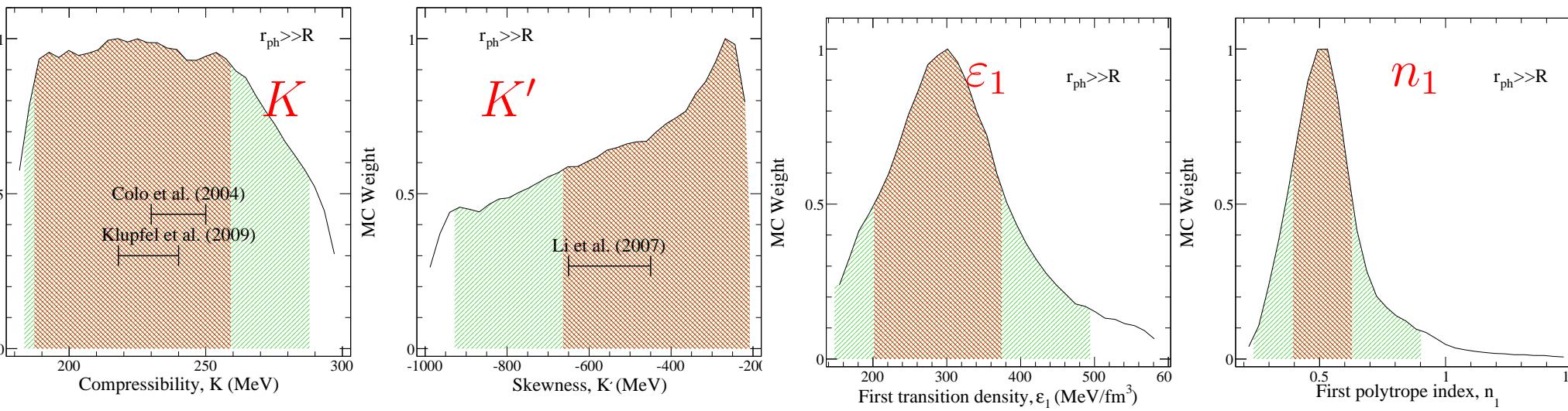
$$P(p_i|D) \propto \int P(D|\mathcal{M}_i) dp_1 dp_2 \dots dp_{i-1} dp_{i+1} \dots dp_{N_p} dM_1 dM_2 \dots dM_{N_M}$$

Bayesian TOV Inversion

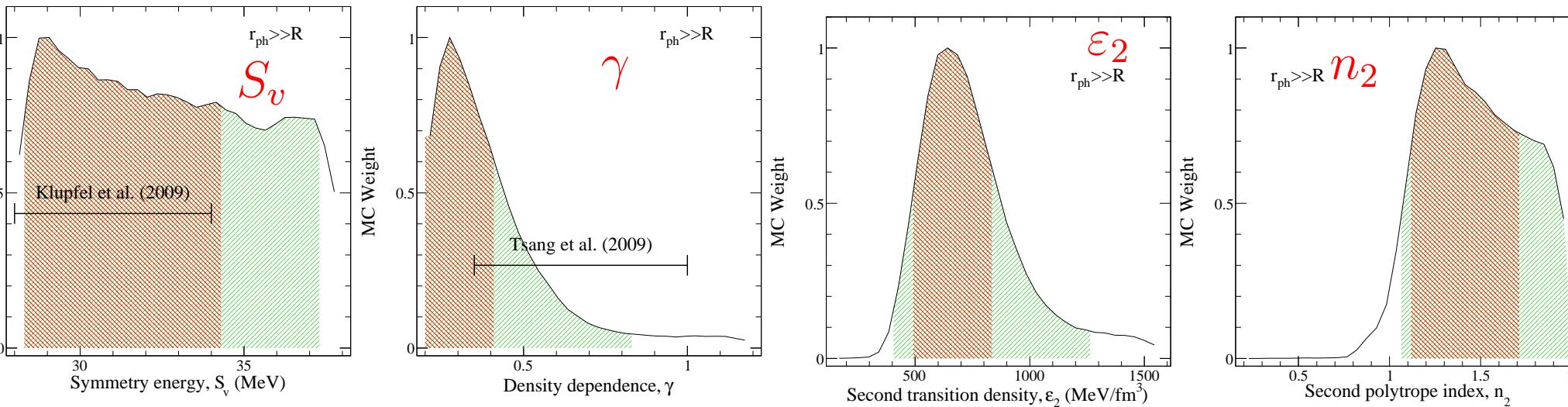
- $\varepsilon < 0.5\varepsilon_0$: EOS from BBP and NV
- $0.5\varepsilon_0 < \varepsilon < \varepsilon_1$: EOS parametrized by K, K', S_v, γ
- $\varepsilon_1 < \varepsilon < \varepsilon_2$: EOS is polytrope with n_1 ; $\varepsilon > \varepsilon_2$: EOS is polytrope with n_2
- A-priori EOS parameters ($K, K', S_v, \gamma, \varepsilon_1, n_1, \varepsilon_2, n_2$) uniformly distributed
- M and R probability distributions for 7 neutron stars treated equally
($0.8 M_\odot < M < 2.5 M_\odot$; $5 \text{ km} < R < 18 \text{ km}$)



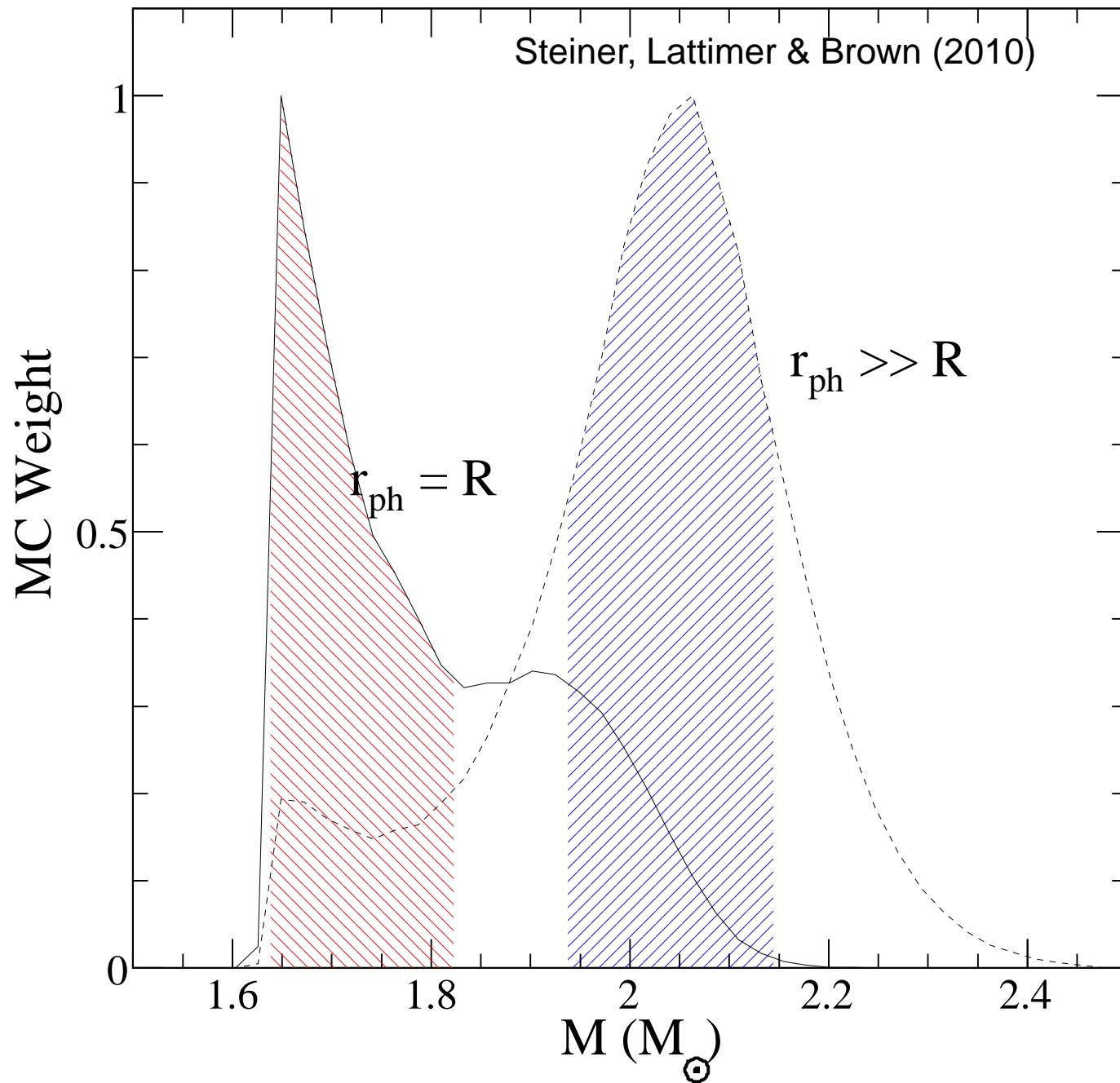
Inferred Model EOS Parameters



Steiner, Lattimer & Brown (2010)



Maximum Mass Probability Distributions



Conclusions

- In spite of large estimated errors for individual sources, when taken as an ensemble they predict a remarkably tight pressure-density relation.
- Nuclear symmetry energy is predicted to be very soft, $\gamma \simeq 0.3 \pm 0.1$, consistent with pure neutron matter results ($\gamma \simeq 0.3 \pm 0.3$). Incompatible with Shen et al. EOS.
- Estimated radii of $1.4 M_{\odot}$ stars are $\simeq 11.3 \pm 0.3$ km.
- The maximum neutron star mass is estimated to be near or above $2 M_{\odot}$.
- The estimated maximum mass and stiffness of the high-density EOS are sensitive to interpretations of photospheric radius expansion bursts.