

# Clusters in Dense Matter and the Equation of State

**Stefan Typel**

Excellence Cluster 'Universe', Technische Universität München  
GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt

in collaboration with

**Gerd Röpke** (Universität Rostock)

**Thomas Klähn** (Uniwersytet Wrocławski)

**David Blaschke** (Uniwersytet Wrocławski)

**Hermann Wolter** (LMU München)

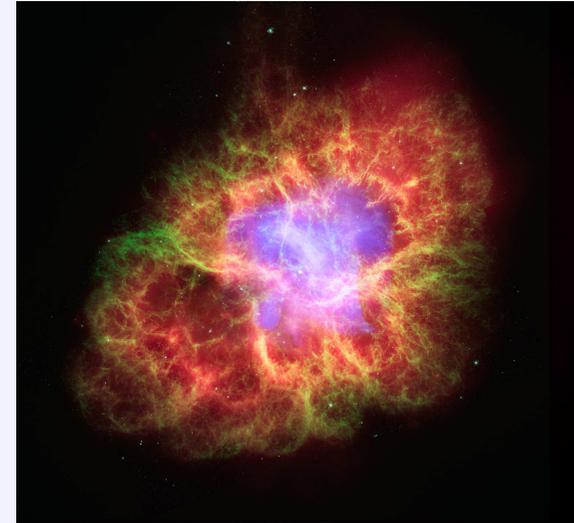
**Maria Voskresenskaya** (GSI Darmstadt)

**11th Symposium on Nuclei in the Cosmos**

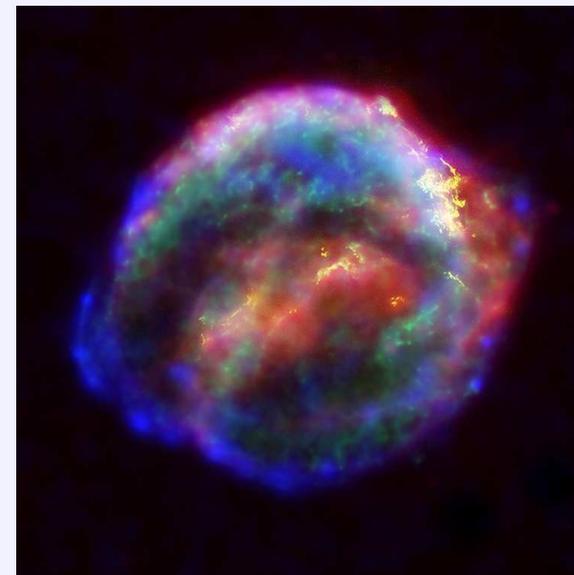
# The Life and Death of Stars

when the era of nuclear fusion reactions ends:

- last phases in the life of a massive star  
( $8M_{\text{sun}} \lesssim M_{\text{star}} \lesssim 30M_{\text{sun}}$ )  
⇒ core-collapse supernova  
⇒ neutron star or black hole



X-ray: NASA/CXC/J.Hester (ASU)  
Optical: NASA/ESA/J.Hester & A.Loll (ASU)  
Infrared: NASA/JPL-Caltech/R.Gehrz (Univ. Minn.)

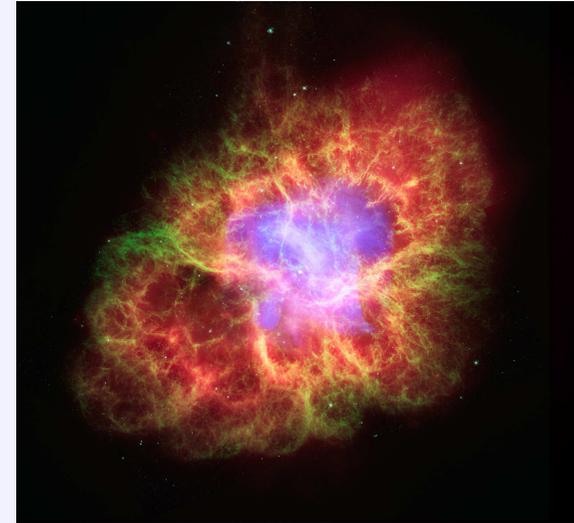


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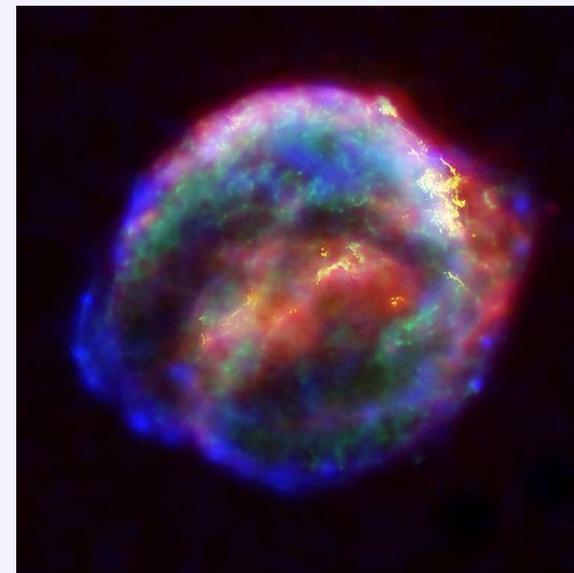
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⇒ core-collapse supernova  
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- essential ingredient in  
astrophysical model calculations:  
**equation of state (EoS) of dense matter**  
⇒ dynamical evolution of supernova  
⇒ static properties of neutron star  
⇒ conditions for nucleosynthesis  
⇒ energetics, chemical composition,  
transport properties, . . .



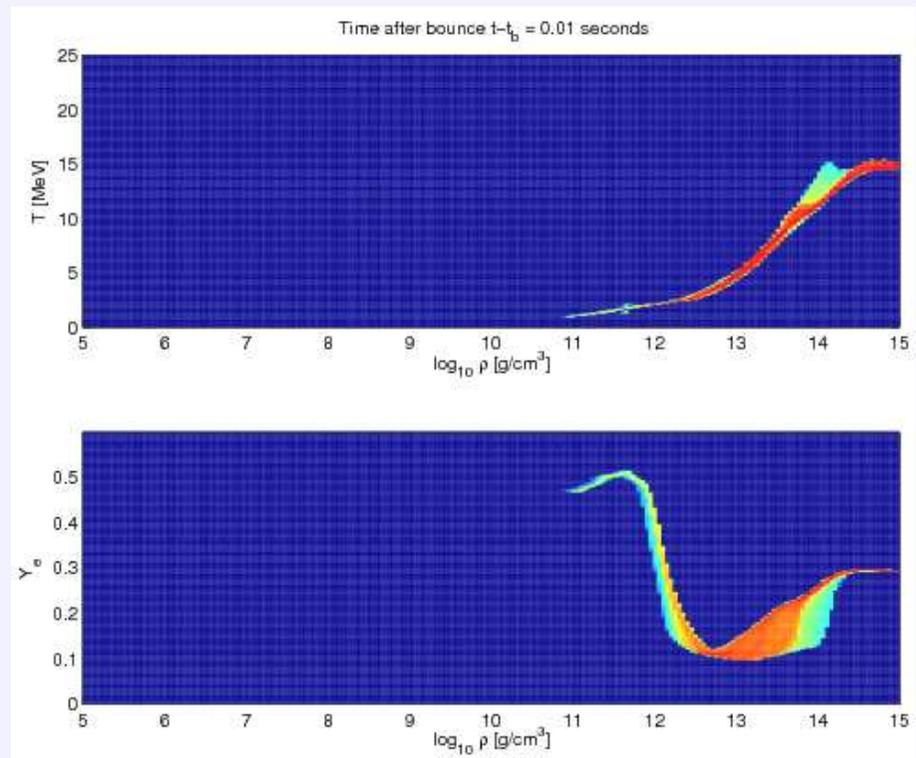
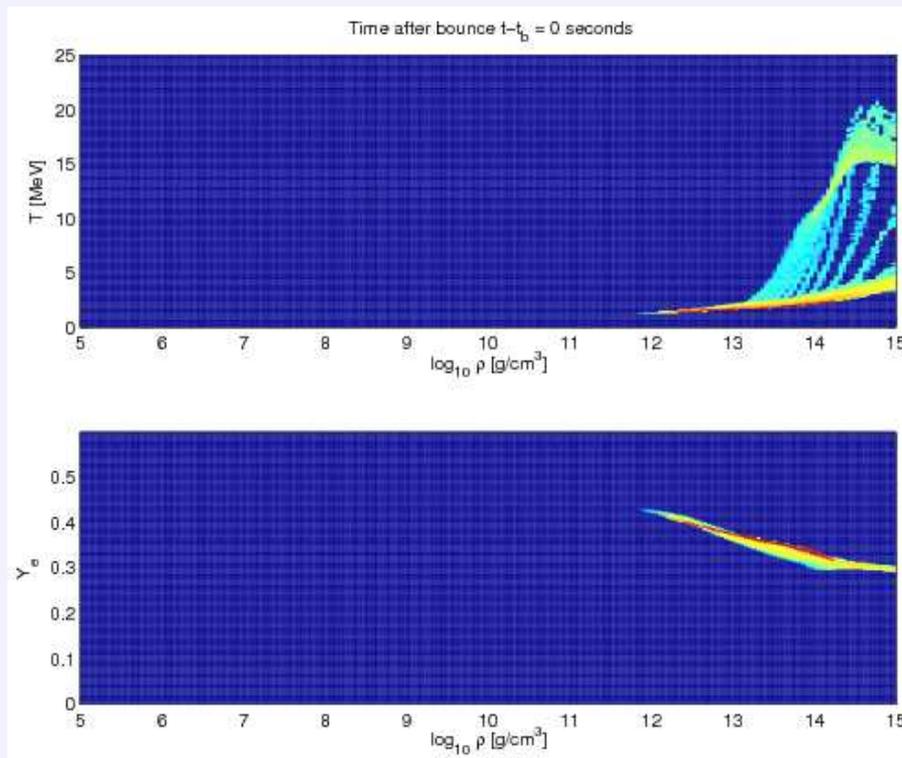
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# Thermodynamical Conditions

- **densities:**  $10^{-9} \lesssim \rho/\rho_{\text{sat}} \lesssim 10$  ( $\rho_{\text{sat}} \approx 2.5 \cdot 10^{14} \text{ g/cm}^3$ )
- **temperatures:**  $0 \text{ MeV} \leq k_B T \lesssim 25 \text{ MeV}$  ( $\hat{=} 2.9 \cdot 10^{11} \text{ K}$ )
- **electron fraction:**  $0 \leq Y_e \lesssim 0.6$

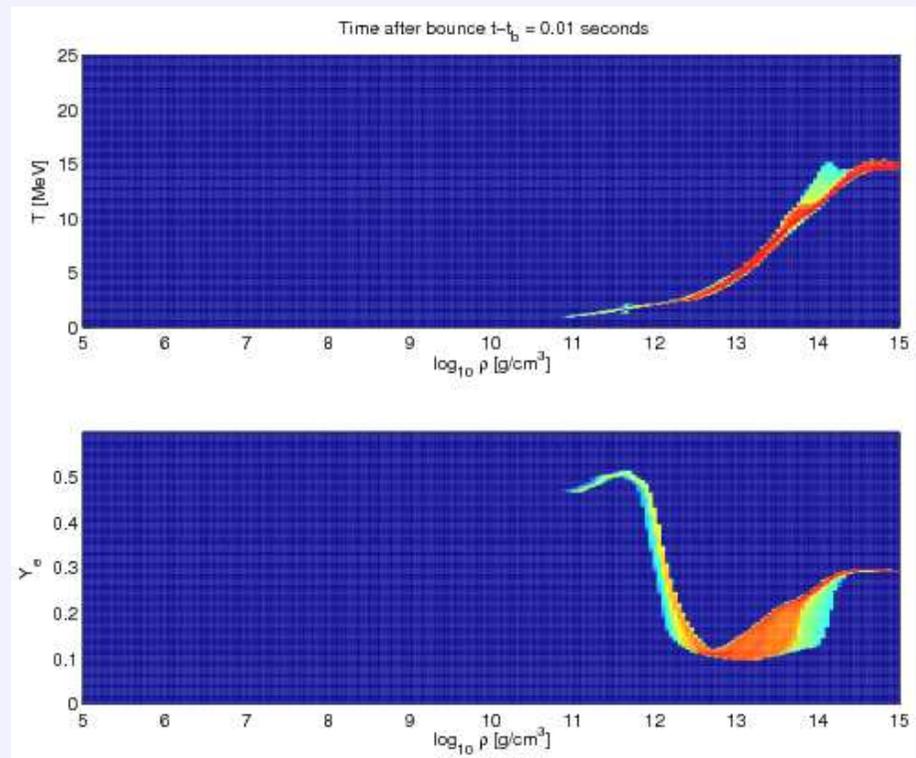
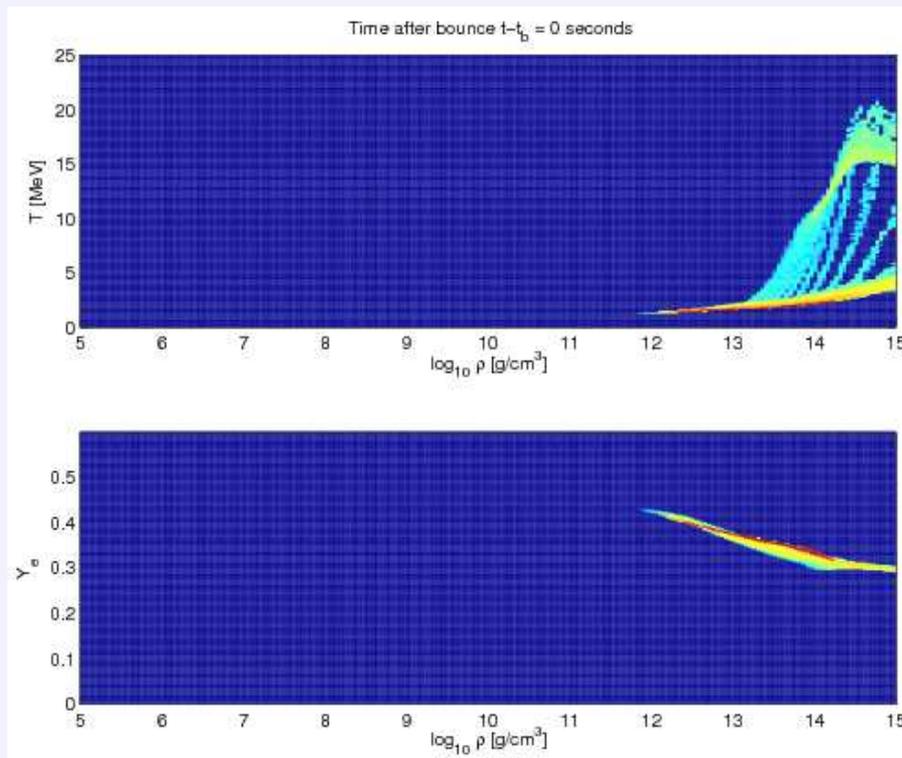


M. Liebendörfer, R. Käppeli, S. Scheidegger, Universität Basel

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⇒ global theoretical description of matter properties is required



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# Equation of State of Dense Matter

- many EoS developed in the past:  
from simple parametrizations to sophisticated models
- many investigations of detailed aspects:  
often restricted to particular conditions (e.g. zero temperature)
- only few EoS used in astrophysical models: most well known
  - J.M. Lattimer, F.D. Swesty (Nucl. Phys. A 535 (1991) 331)
  - H. Shen, H. Toki, K. Oyamatsu, K. Sumiyoshi (Prog. Theor. Phys. 100 (1998) 1013)

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- most difficult problem:  
description of strongly interacting subsystem (hadronic or quark matter)  
in this talk: formation of “clusters” in nuclear matter
- in “standard” astrophysical EoS:  
only nucleons,  $\alpha$  particle and representative heavy nucleus,  
suppression of cluster formation with phenomenological excluded-volume mechanism  
⇒ consider **more microscopic model, more clusters**

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- theoretical models: different points of view
  - **chemical picture:**  
mixture of different nuclear species and nucleons in chemical equilibrium  
problems:
    - properties of constituents independent of medium
    - interaction between particles
    - dissolution of nuclei at high densities

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    - **physical picture:**  
correlations of nucleons  $\Rightarrow$  formation of bound states  
problems:
      - treatment of three-, four-, . . . many-body correlations difficult
      - choice of interaction
- $\Rightarrow$  combination of approaches?

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- **low densities:**
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“liquid-gas” phase transition
  - surface effects and long-range Coulomb interaction
  - inhomogeneous matter
  - formation of “pasta” phases/lattice structures

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**interpolation between low-density and high-density limit needed**

⇒ consider quantum statistical approach and generalized relativistic mean-field model

# Quantum Statistical Approach I

- nonrelativistic finite-temperature Green's function formalism
- starting point: nucleon number densities ( $\tau = p, n$ )

$$n_\tau(T, \tilde{\mu}_p, \tilde{\mu}_n) = 2 \int \frac{d^3k}{(2\pi)^3} \int \frac{d\omega}{2\pi} f_\tau(\omega) S_\tau(\omega) \quad \text{with Fermi distribution } f_\tau(\omega)$$

and spectral function  $S_\tau(\omega)$  depending on self-energy  $\Sigma_\tau$

- expansion of spectral function beyond quasiparticle approximation

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⇒ **generalized Beth-Uhlenbeck description** with

- medium dependent self-energy shifts/binding energies
- generalized scattering phase shifts from in-medium T-matrix

- $T, n_p, n_n \Rightarrow \tilde{\mu}_p, \tilde{\mu}_n \Rightarrow$  free energy  $F(T, n_p, n_n)$  by integration  $\left. \left( \frac{\partial(F/V)}{\partial n_\tau} \right) \right|_{T, n_{\tau'}} = \tilde{\mu}_\tau$

⇒ **thermodynamically consistent** derivation of **EoS**

# Quantum Statistical Approach II

## medium modifications

- single nucleon properties
  - self-energy shift of quasiparticle energy
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# Quantum Statistical Approach II

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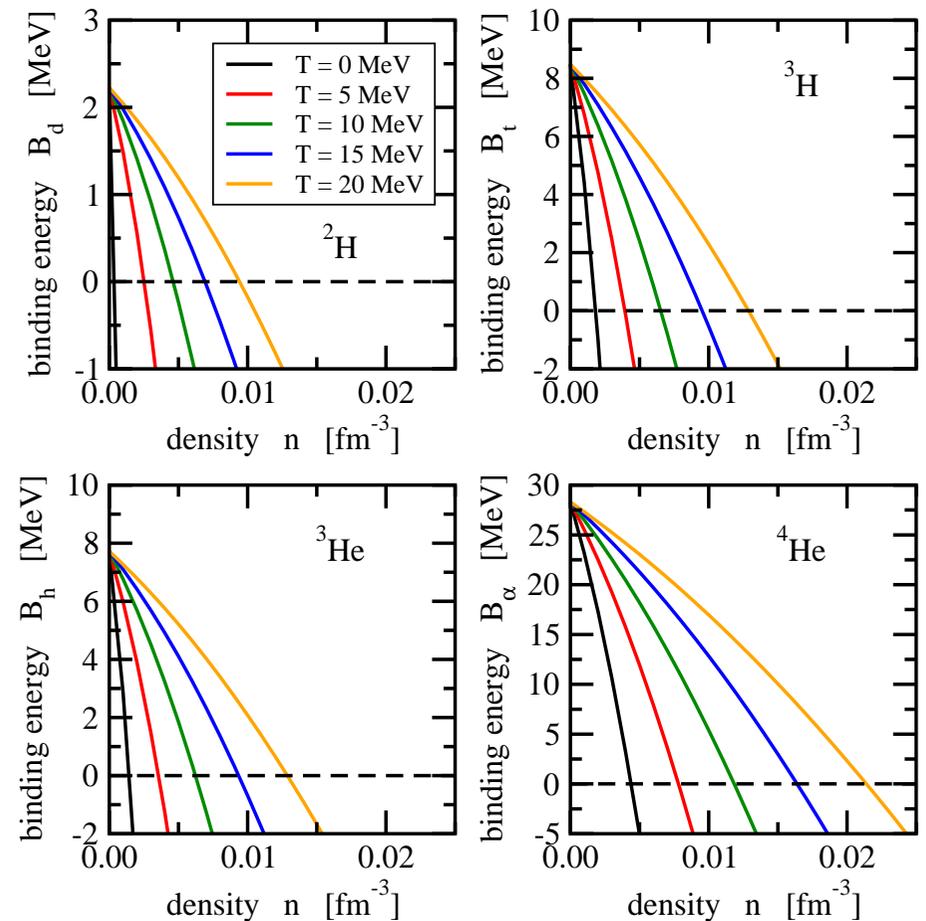
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    - shift of quasiparticle energy from
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(calculation with effective nucleon-nucleon potential)

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(calculation with effective nucleon-nucleon potential)
- ⇒ quasi-particles

symmetric nuclear matter



parametrization used in generalized RMF model

# Generalized Relativistic Mean-Field (RMF) Model

- extended **relativistic Lagrangian density** of Walecka type
  - with **nucleons** ( $\psi_p, \psi_n$ ), **deuterons** ( $\varphi_d^\mu$ ), **tritons** ( $\psi_t$ ), **helions** ( $\psi_h$ ),  **$\alpha$ -particles** ( $\varphi_\alpha$ ), **mesons** ( $\sigma, \omega_\mu, \vec{\rho}_\mu$ ), **electrons** ( $\psi_e$ ) and **photons** ( $A_\mu$ ) as degrees of freedom
    - only **minimal** (linear) **meson-nucleon couplings**
    - **density-dependent** meson-nucleon **couplings**  $\Gamma_i$

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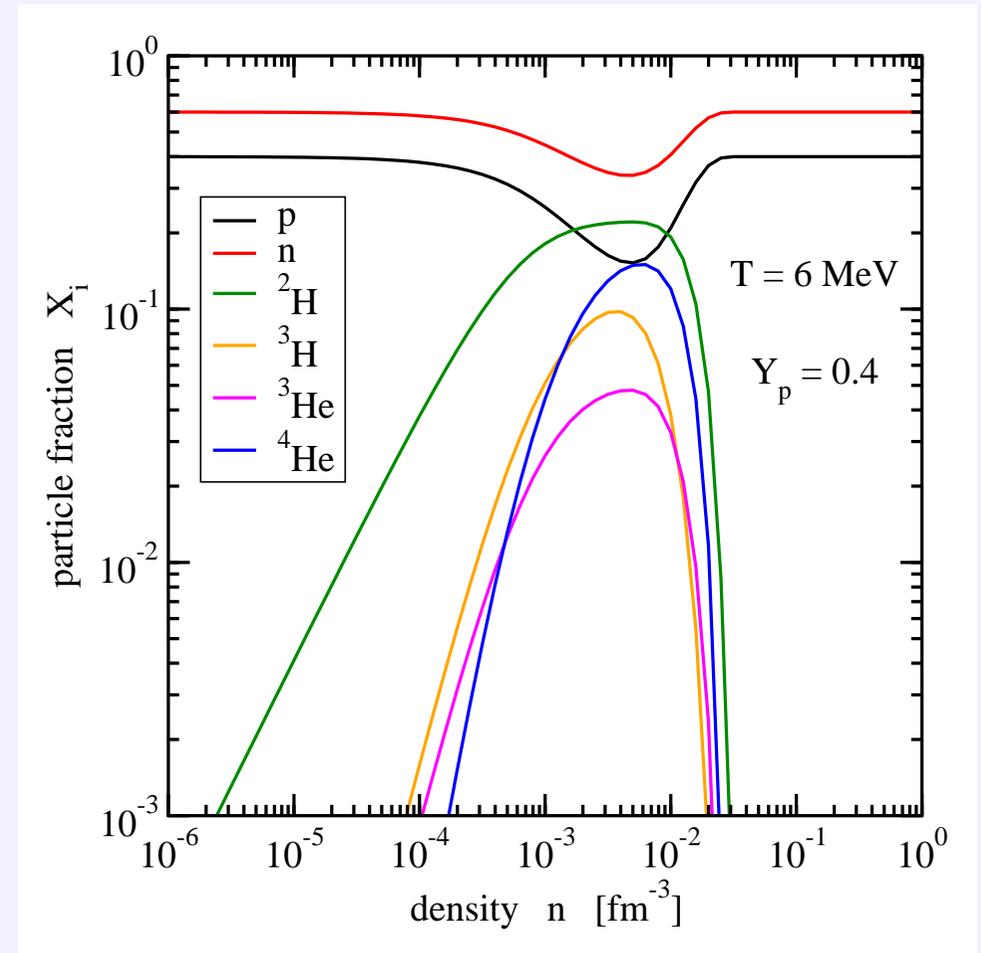
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  - ⇒ nucleon/cluster/meson/photon **field equations**,  
solved selfconsistently in **mean-field approximation**  
(classical meson/photon fields, Hartree approximation, no-sea approximation)

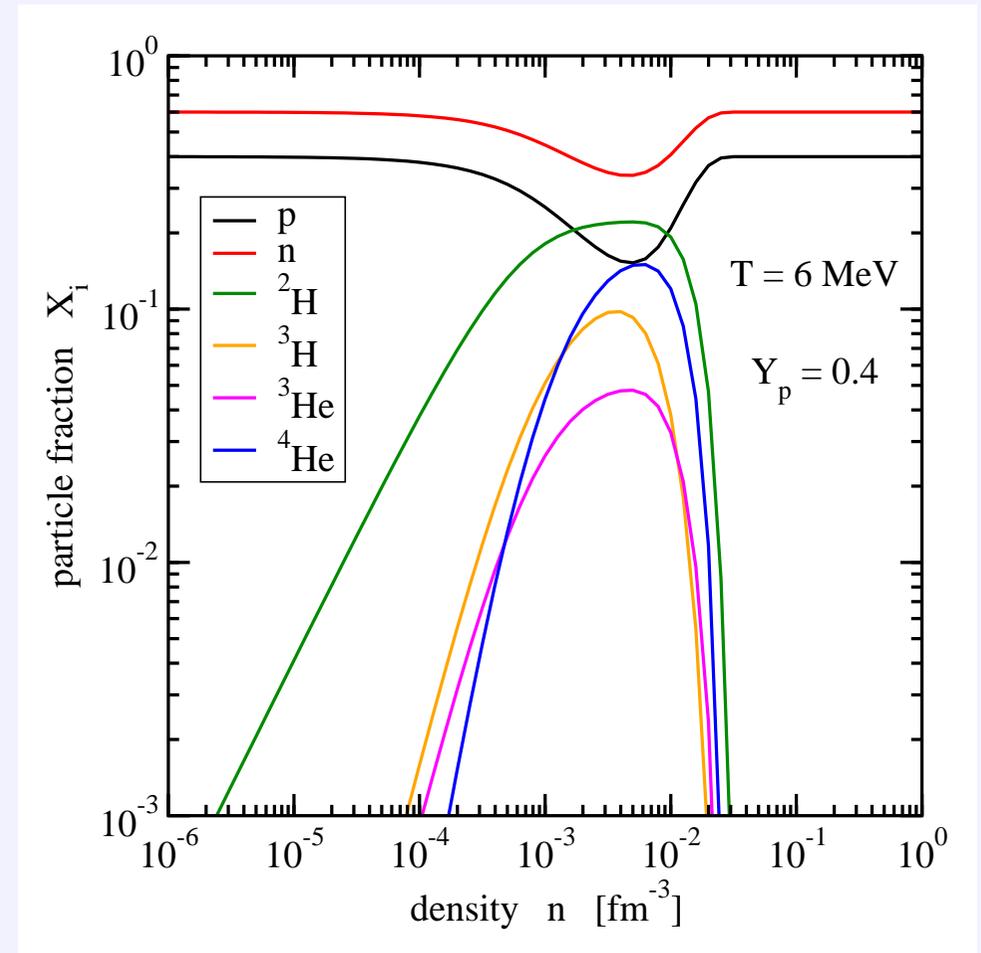
# EoS with Light Clusters - Generalized RMF Model

- consider 2-, 3-, and 4-body correlations in the medium
  - presently only bound states (deuterons, tritons, helions, and alphas)
  - scattering contributions neglected so far
- Mott effect: clusters dissolve at high densities
- correct limits at low and high densities



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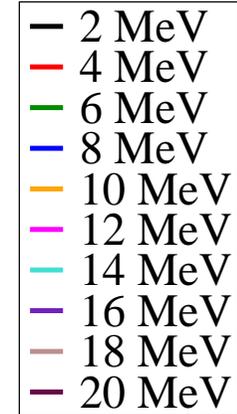
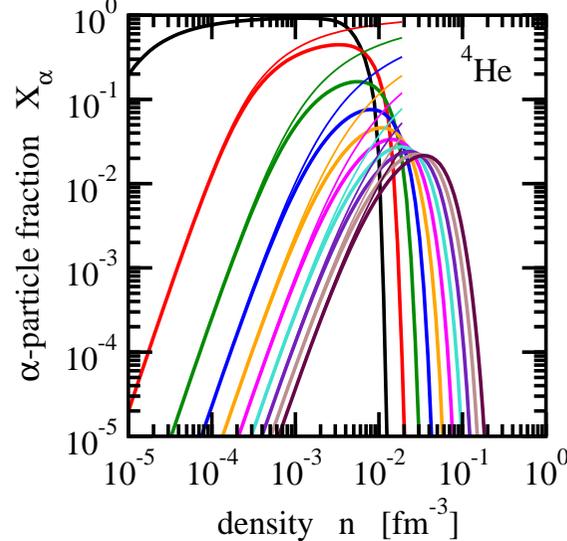
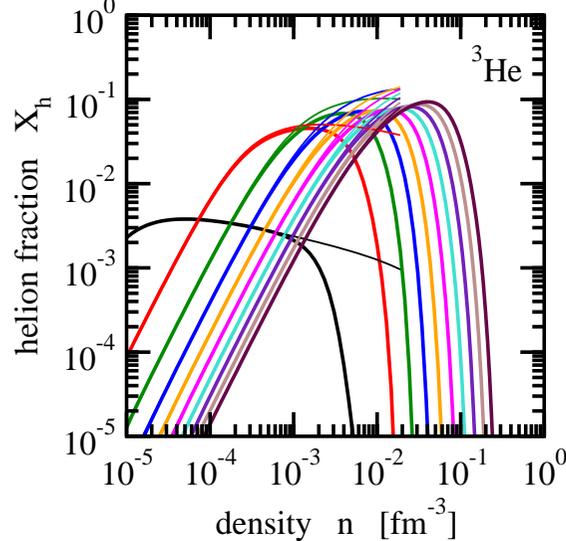
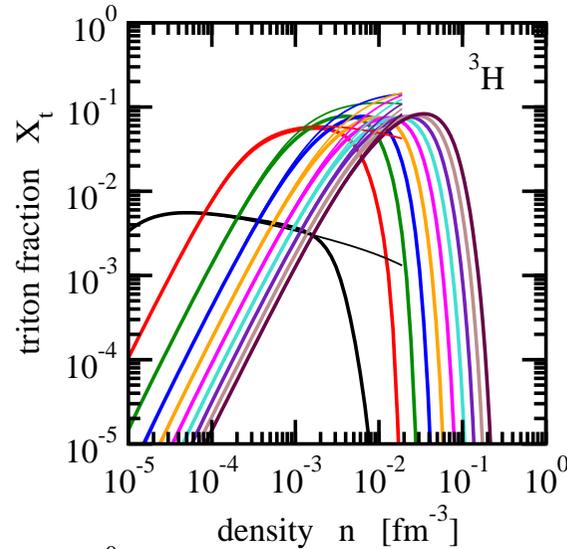
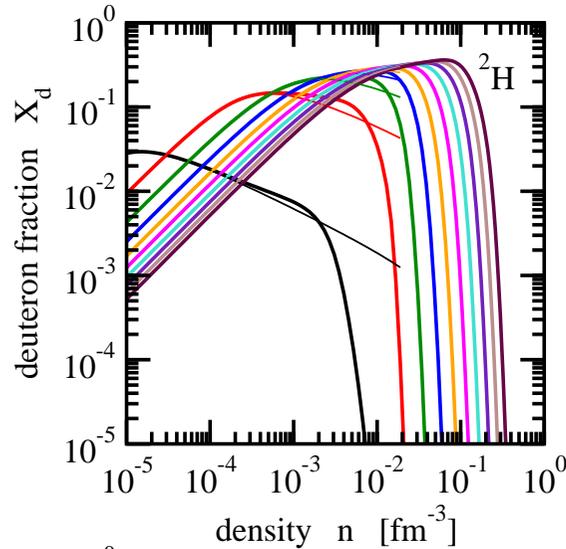
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- **Mott effect**: clusters dissolve at high densities
- **correct limits** at low and high densities
- no heavy clusters/phase transition included here
- **medium dependence** of couplings and binding energies
  - ⇒ “**rearrangement**” contributions in self-energies and source densities essential for **thermodynamical consistency**



# EoS with Light Clusters - Cluster Fractions

symmetric nuclear matter

generalized RMF model vs. NSE (thin lines)

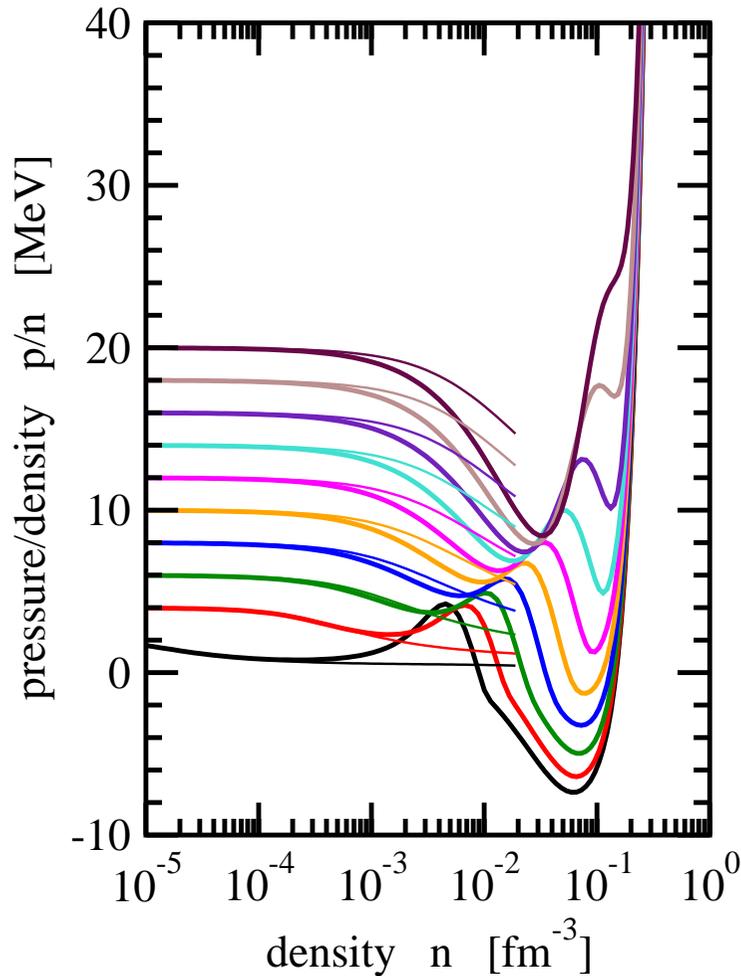


# EoS with Light Clusters - Pressure/Density

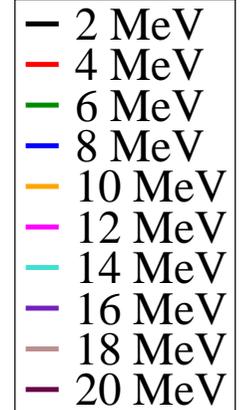
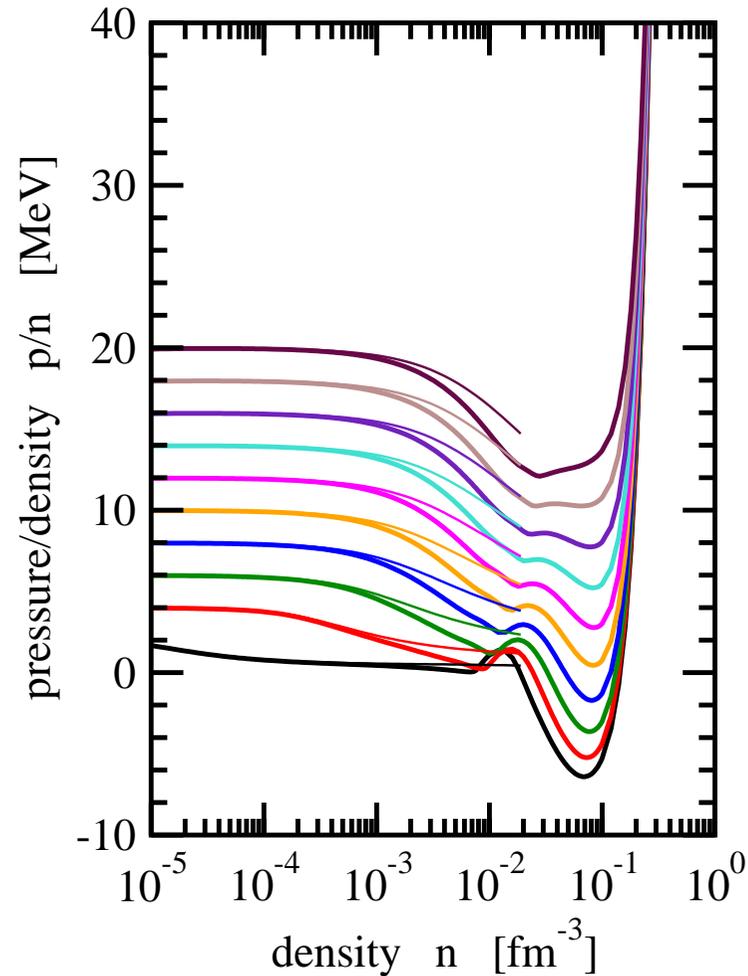
symmetric nuclear matter

$$\lim_{n \rightarrow 0} (p/n) = T \quad (\text{ideal gas})$$

generalized RMF vs. NSE



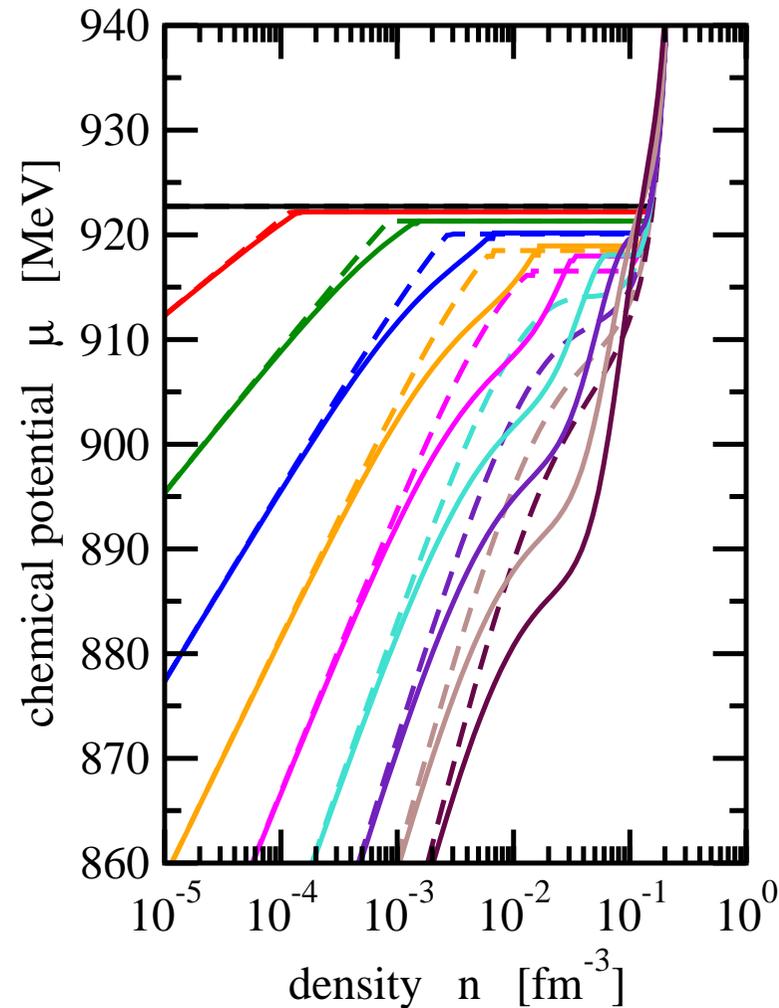
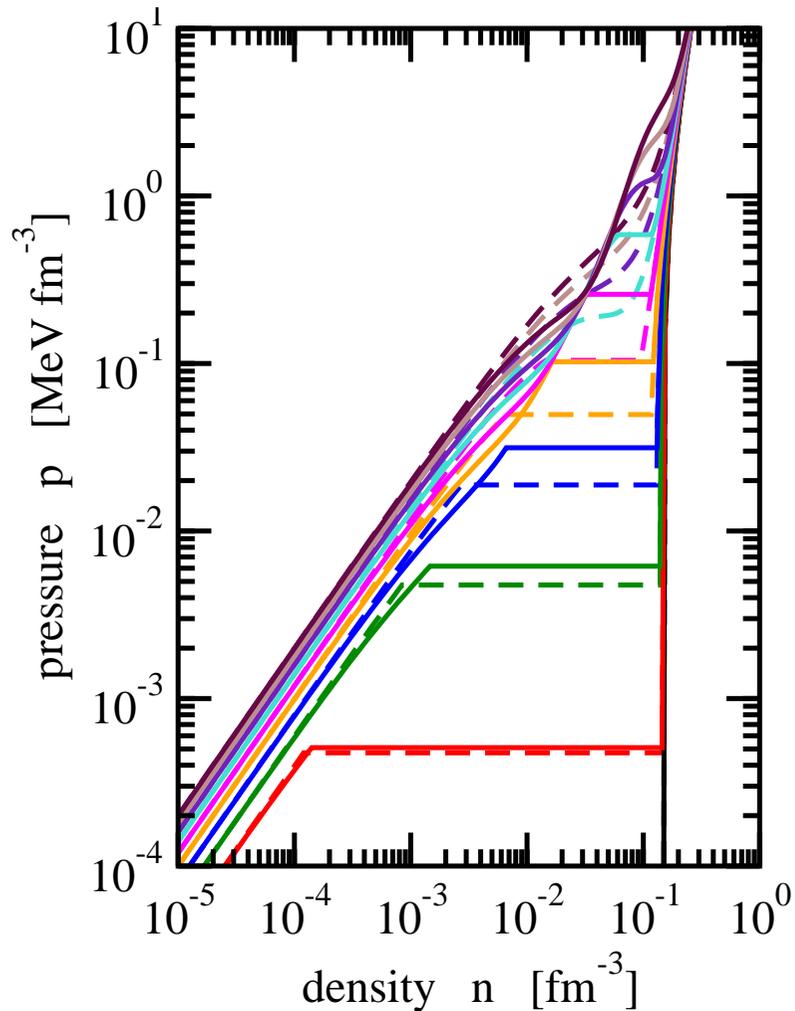
QS approach vs. NSE



# Phase Transition - Pressure and Chemical Potential

symmetric nuclear matter (Maxwell construction sufficient)

RMF model without (dashed lines) and with (solid lines) clusters

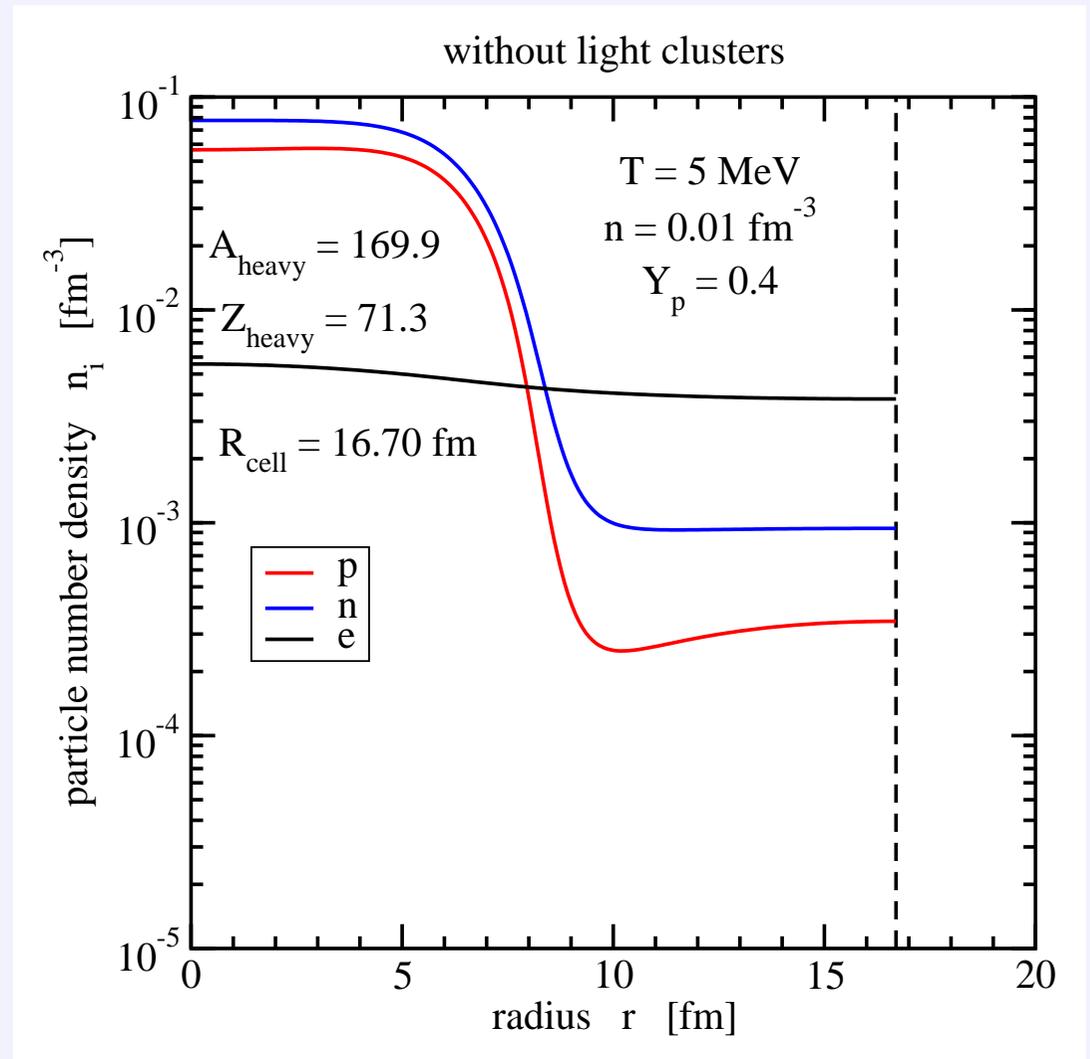


# Heavy Clusters

- liquid-gas phase transition:  
separation of low-/high-density phases,  
no surface or Coulomb effects

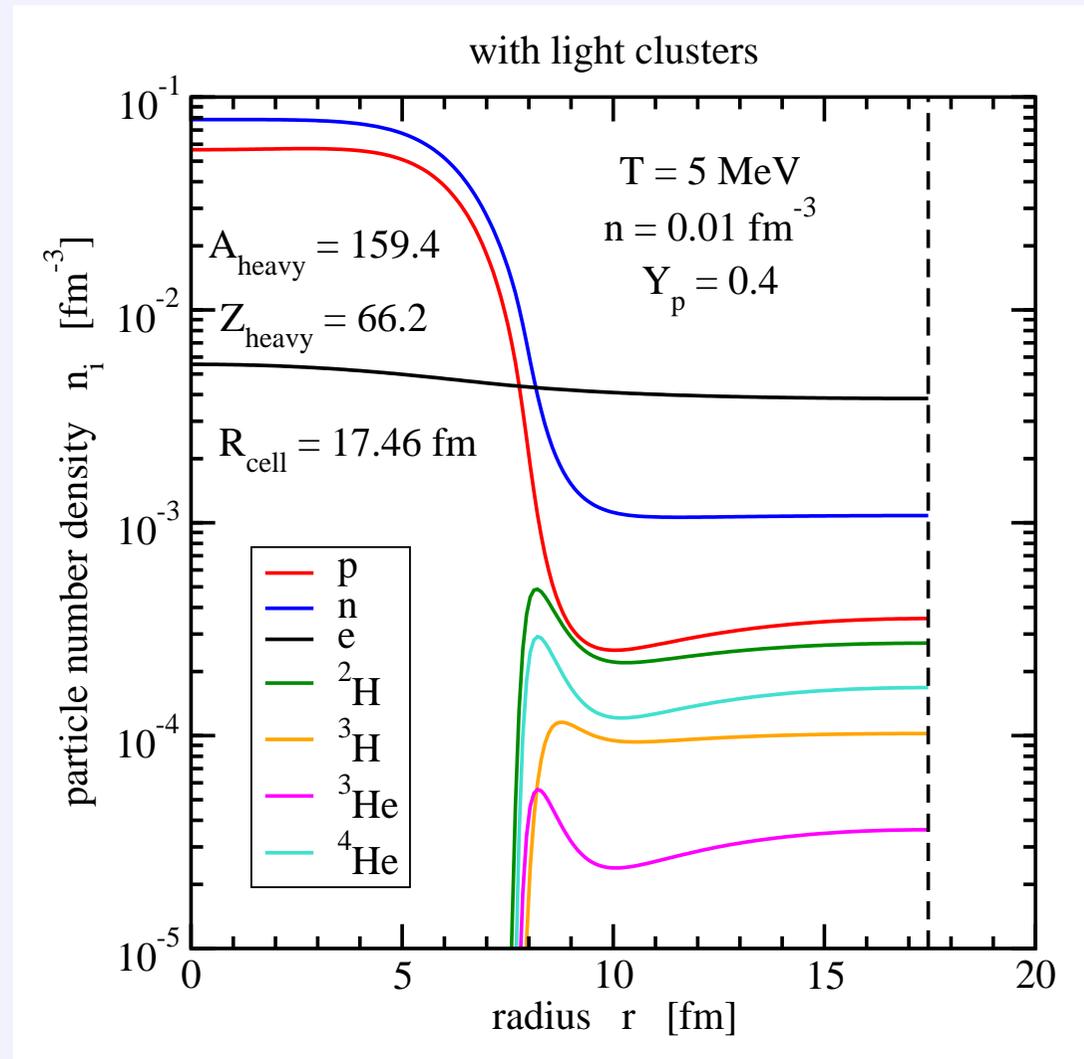
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- first step in improvement:  
spherical Wigner-Seitz cell calculation
  - generalized RMF model
  - Thomas-Fermi approximation
  - electrons for charge compensation
  - heavy nucleus surrounded by  
gas of nucleons and light clusters



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gas of nucleons and light clusters
- first self-consistent calculation with  
interacting nucleons, light clusters  
and electrons



# Symmetry Energy I

- **general definition** for zero temperature:

$$E_s(n) = \frac{1}{2} \frac{\partial^2 E}{\partial \beta^2} \frac{1}{A}(n, \beta) \Big|_{\beta=0} \quad \beta = \frac{n_n - n_p}{n_n + n_p}$$

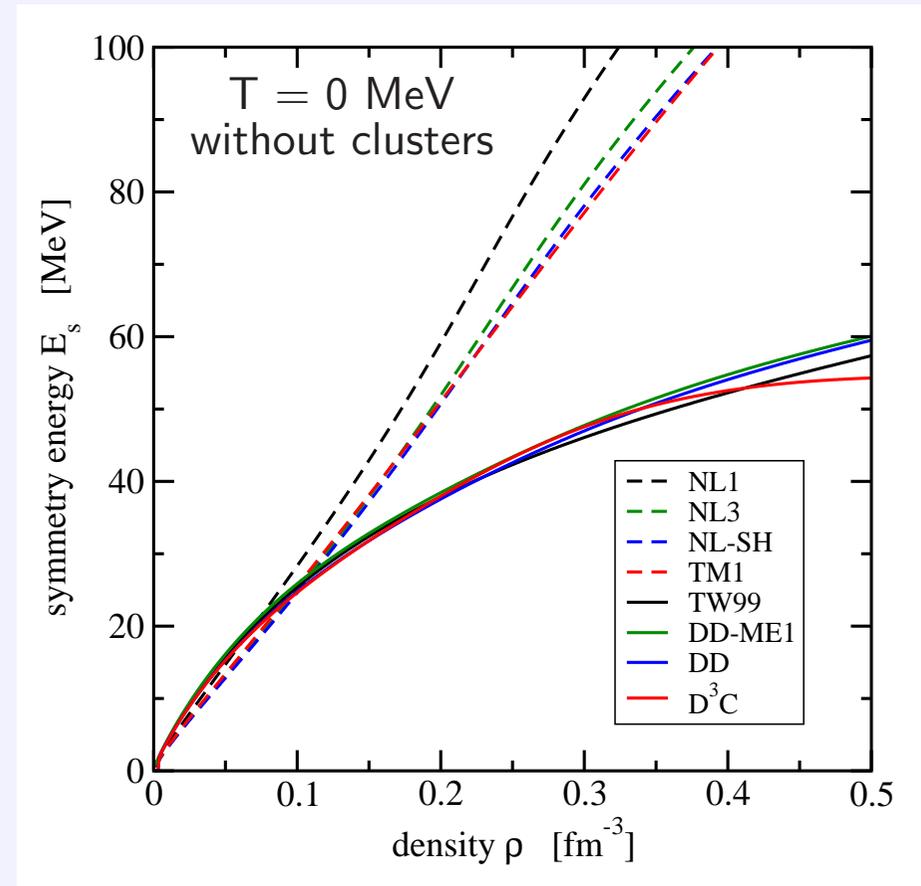
⇒ nuclear matter parameters

$$J = E_s(n_{\text{sat}}) \quad L = 3n \frac{d}{dn} E_s \Big|_{n=n_{\text{sat}}}$$

- **correlation**: neutron skin thickness  
⇔ slope of neutron matter EoS (⇔  $L$ )

B. A. Brown, Phys. Rev. Lett. 85 (2000) 5296,

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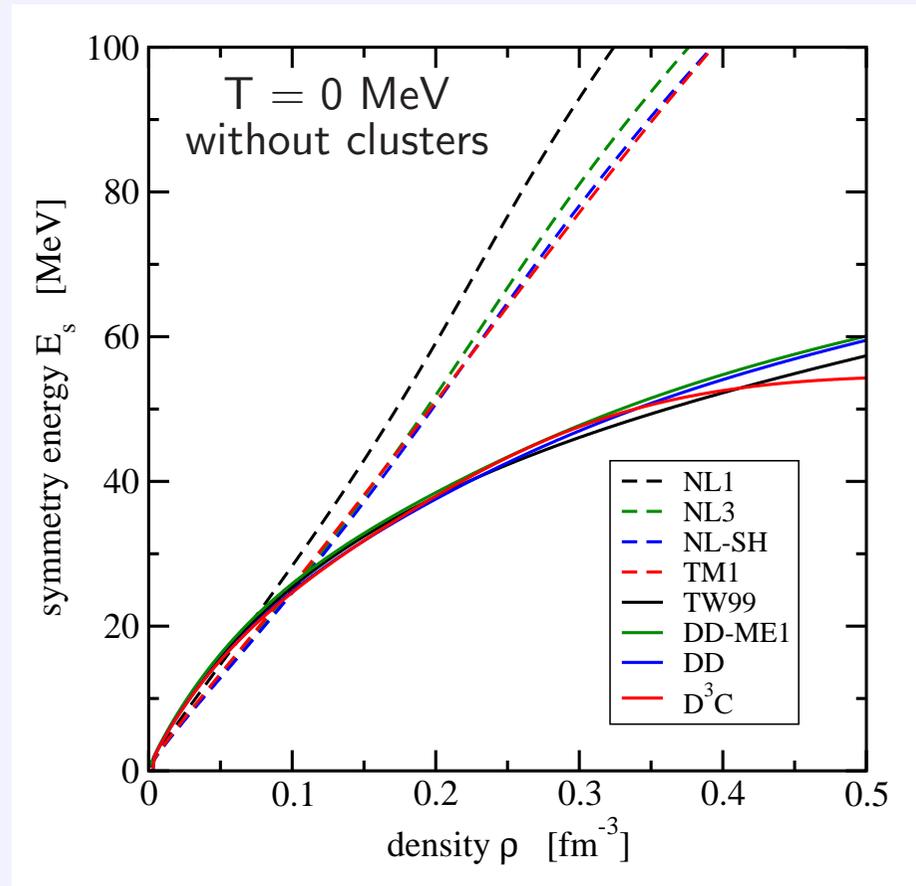
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- **with clusters and at finite temperatures:**
  - use finite differences

$$E_{\text{sym}}(n) = \frac{1}{2} \left[ \frac{E}{A}(n, 1) - 2\frac{E}{A}(n, 0) + \frac{E}{A}(n, -1) \right]$$

**effects of cluster formation? experimental observation?**



# Symmetry Energy II

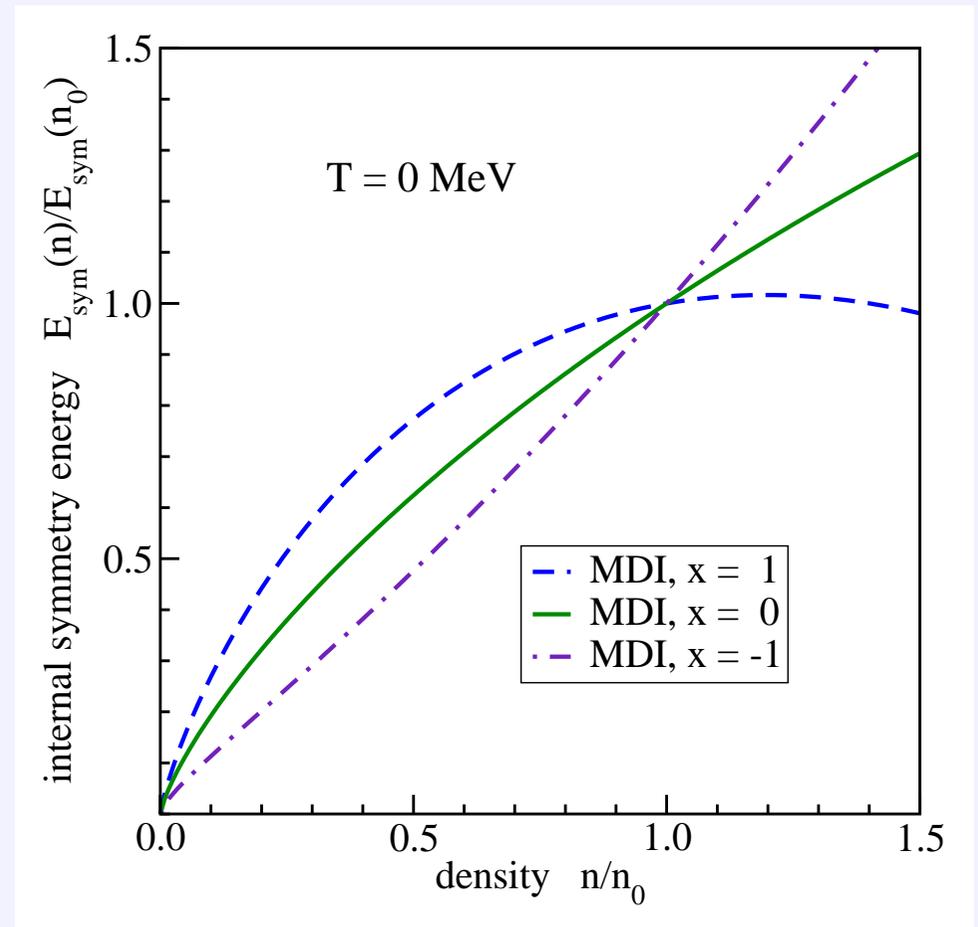
temperature  $T = 0$  MeV

- mean-field models without clusters

e.g. model with momentum-dependent interaction (MDI), parameter  $x$  controls density dependence of  $E_{\text{sym}}$

(B. A. Li et al., Phys. Rep. 464 (2008) 113)

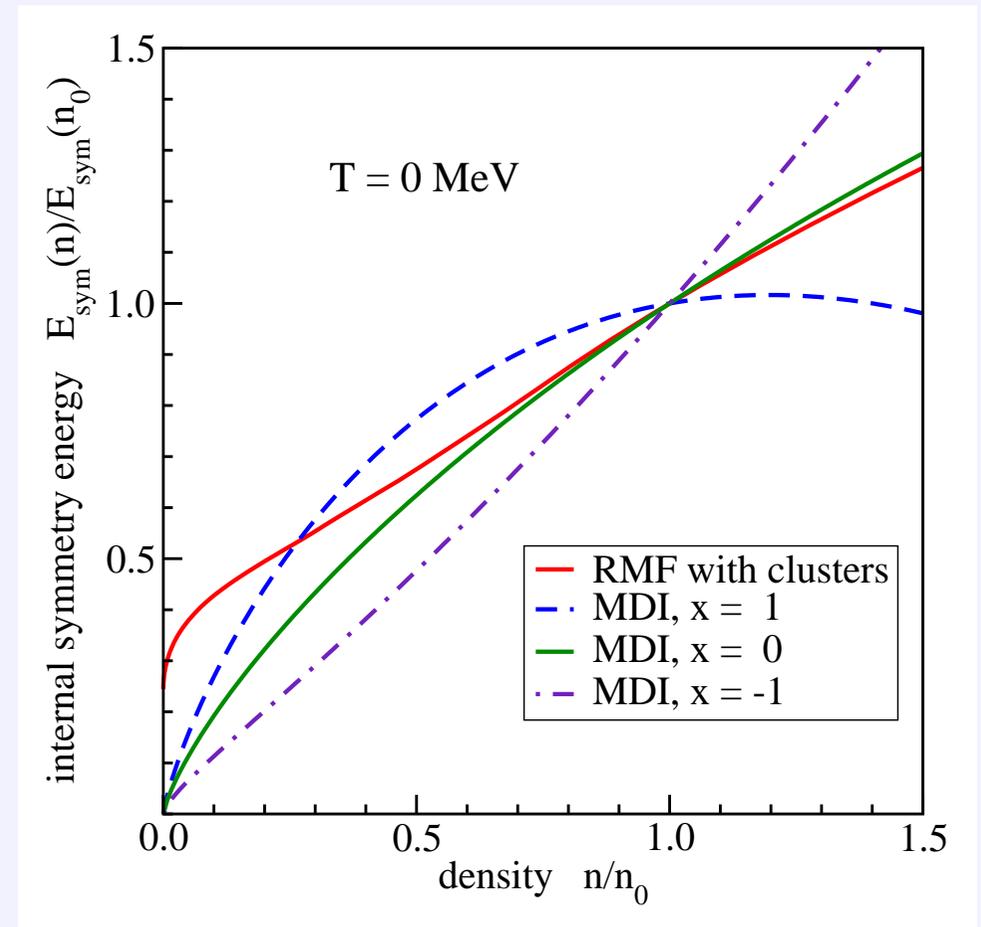
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  - ⇒ low-density behaviour not correct
- RMF model with (heavy) clusters
  - ⇒ increase of  $E_{\text{sym}}$  at low densities due to formation of clusters
  - ⇒ finite symmetry energy in the limit  $n \rightarrow 0$



# Symmetry Energy III

## finite temperature

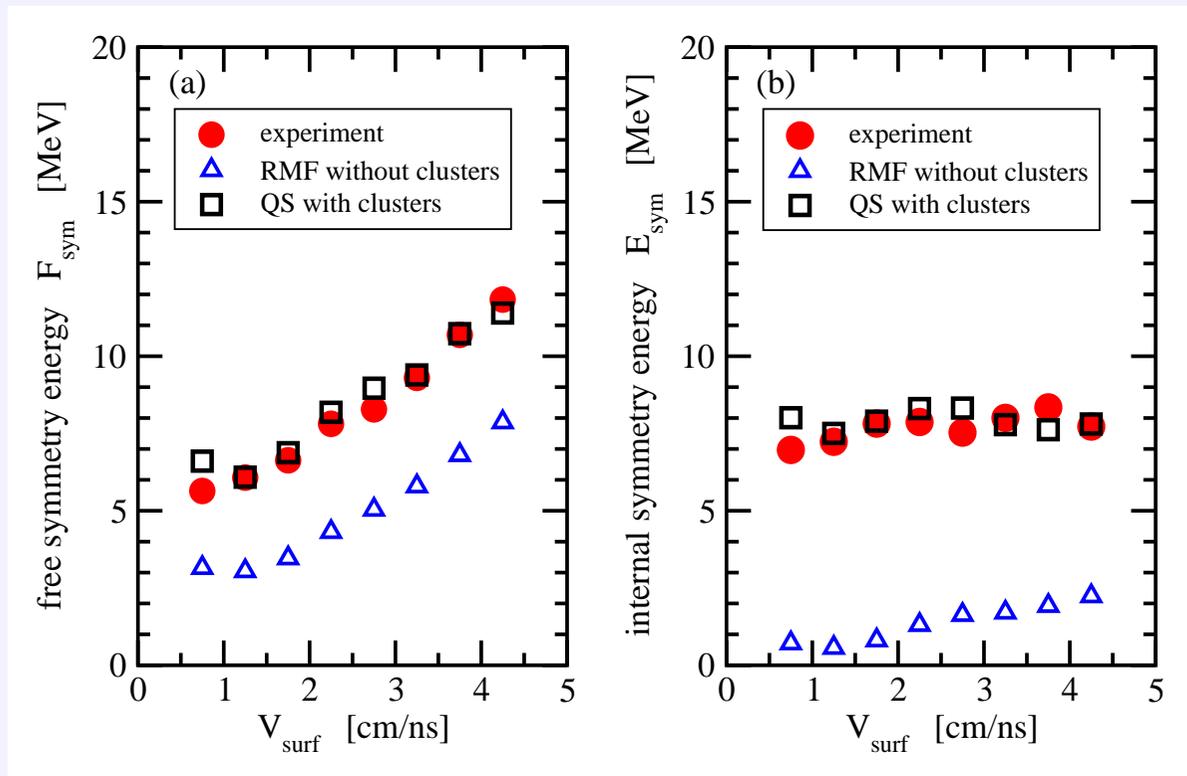
- experimental determination of symmetry energy
  - heavy-ion collisions of  $^{64}\text{Zn}$  on  $^{92}\text{Mo}$  and  $^{197}\text{Au}$  at 35  $A$  MeV  
temperature, density, free symmetry energy derived as functions of  
parameter  $v_{\text{surf}}$  (measures time when particles leave the source)  
(S. Kowalski et al., Phys. Rev. C 75 (2007) 014601)

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- symmetry energies in RMF calculation without clusters are too small
- very good agreement with QS calculation with light clusters



# Summary and Outlook

- **theoretical models of EoS with clusters**
  - quantum statistical approach (QS)
  - generalized relativistic mean-field model (gRMF)
  - both thermodynamically consistent
  - correct limits at low and high densities
  - difference in details

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- **nuclear matter at low densities**
  - formation of clusters with medium dependent properties
  - modification of thermodynamical properties/symmetry energies
  - change of phase transition boundaries

for details see Phys. Rev. C **81**, 015803 (2010) and Phys. Rev. Lett. **104**, 202501 (2010)

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- **future**
    - further improvement of RMF parametrization (low-density limit)
    - application to astrophysical models
- ⇒ CompStar (compstar-esf.org) initiative:  
repository of modern EoS for astrophysical applications